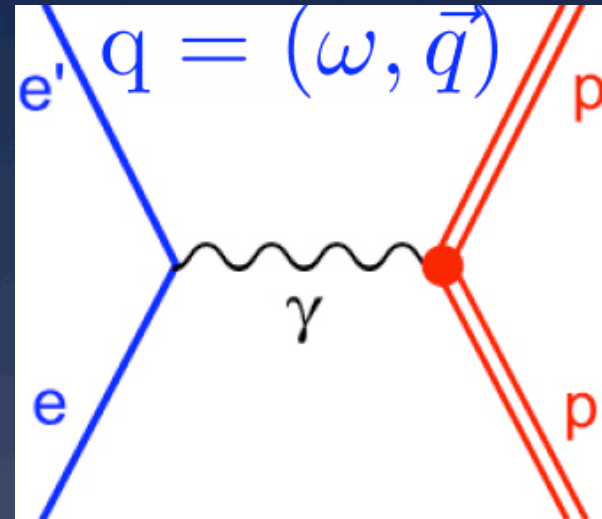


Nucleon Form Factor Experiments: Present and Future

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University of Virginia

Elastic Electro-Magnetic Form Factors

The hadronic current in one photon exchange



$$J^\mu = e\bar{u}(p') [F_1(q^2)\gamma^\mu + i\frac{\kappa}{2M}q_\nu\sigma^{\mu\nu}F_2(q^2)]u(p)$$

F_1 , Dirac form factor; helicity conserving

F_2 , Pauli form factor; helicity non-conserving

Or in terms of Sachs form factors

$$G_E = F_1 - \kappa\tau F_2$$

$$G_M = F_1 + \kappa F_2, \tau = \frac{Q^2}{4M^2}$$

So the elastic cross section

$$\frac{d\sigma}{d\Omega} = \frac{d\sigma}{d\Omega}\Bigg|_{\text{Mott}} \frac{E'}{E} \left[\frac{G_E^2 + \tau G_M^2}{1 + \tau} + 2\tau G_M^2 \tan^2 \frac{\theta}{2} \right]$$

Elastic Electro-Magnetic Form Factors

- Elastic form factors parameterize the properties of the quark and gluon many body system of the nucleon
- They provide excellent testing ground for QCD-inspired models

In the Breit Frame, G_E and G_M are the Fourier transforms of charge and magnetization distributions of the nucleon

$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega} \right)_M \left\{ G_E^2 + \frac{\tau}{\epsilon} G_M^2 \right\} / (1 + \tau)$$

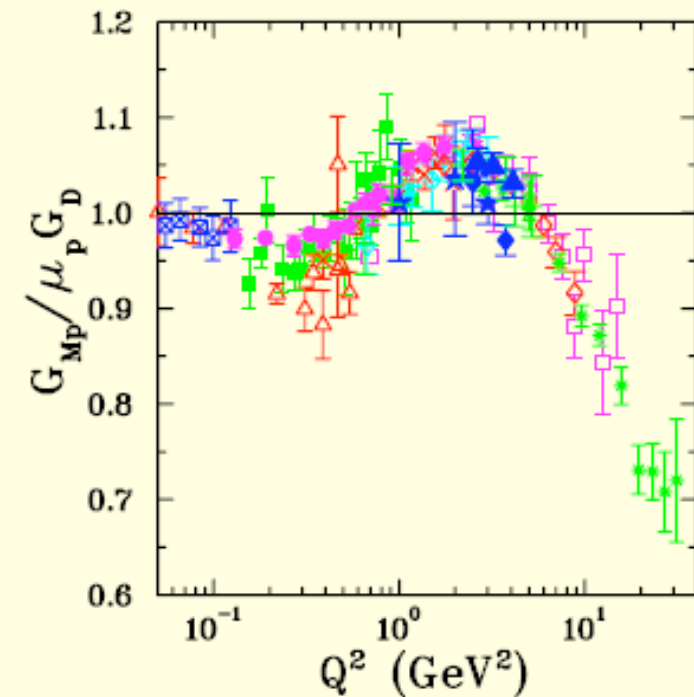
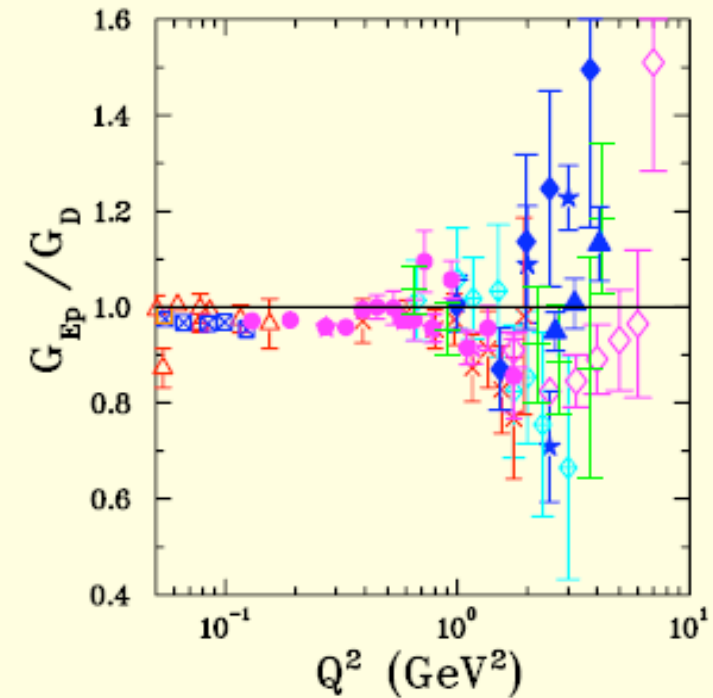
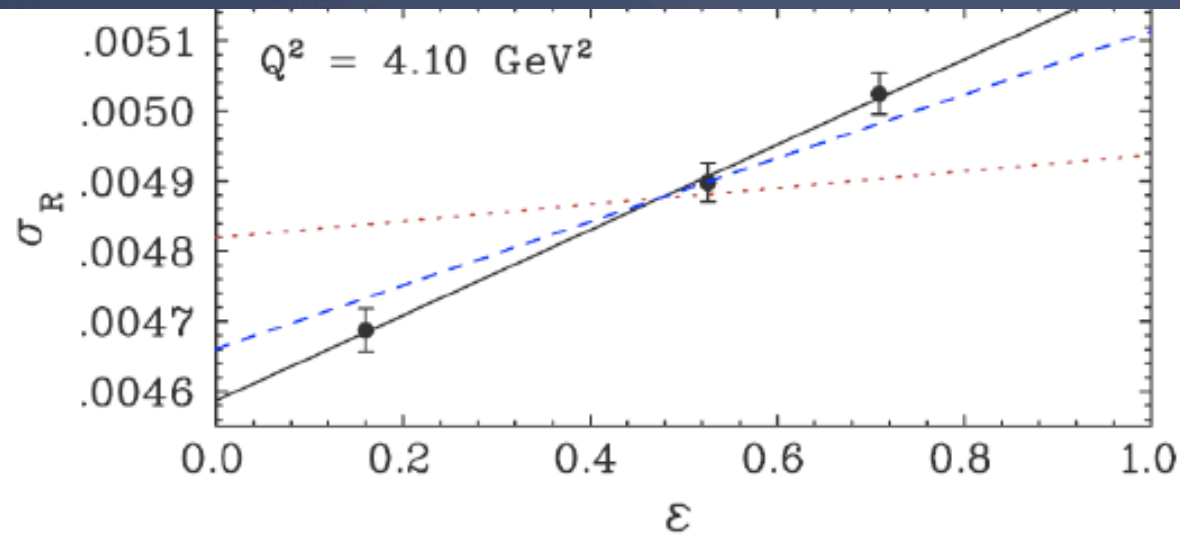
$$\tau = \frac{Q^2}{4M^2}$$

$$\epsilon = \frac{1}{1 + 2(1 + \tau) \tan^2(\frac{\theta_s}{2})}$$

$$\begin{aligned} \sigma_R &= [\epsilon(1+\tau)/\tau] [\sigma_{\text{exp}}/\sigma_{\text{Mott}}] = \\ &= G_{Mp}^2 + \epsilon G_{Ep}^2/\tau \end{aligned}$$

Separation of Electric and Magnetic Form Factors using cross-section data

$$\begin{aligned}\sigma_R &= [\varepsilon(1+\tau)/\tau] [\sigma_{\text{exp}}/\sigma_{\text{Mott}}] = \\ &= G_{Mp}^2 + \varepsilon G_{Ep}^2/\tau\end{aligned}$$



Polarization method to measure GE/GM

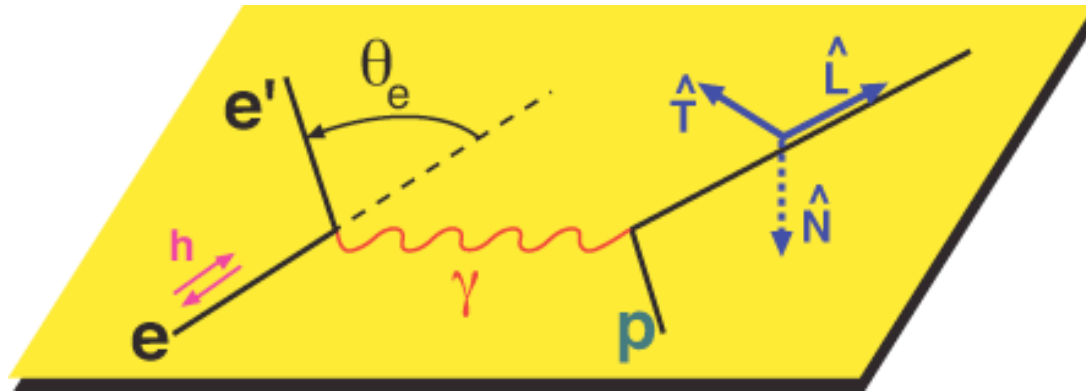
accuracy of form-factor measurements can be significantly improved by measuring the beam helicity asymmetry with a polarized target or with recoil polarimetry

Akhiezer et al., *Sov. Phys. JETP* 6, 588 (1958) and
Arnold, Carlson and Gross, *PR C* 23, 363 (1981)

Finally became possible with Jefferson lab, Mainz and Bates:

- Polarized beam with high intensity ($\sim 100 \mu A$) and high polarization ($>70\%$)
- Beam polarimeters with 1-3 % absolute accuracy
- Polarized targets with a high polarization or
- Recoil polarimeters with large analyzing powers

Recoil polarization method to measure GE_p/GMp



- Ratio cancels beam polarization and analyzing power
- GE/GM measured from a single measurement: no errors due to ε dependant effects
- Not sensitive to: spectrometer solid angle, target density, trigger and detector inefficiencies, beam charge asymmetry, false asymmetries in FPP
- Main sensitivity is to spin transport

$$P_n = 0$$

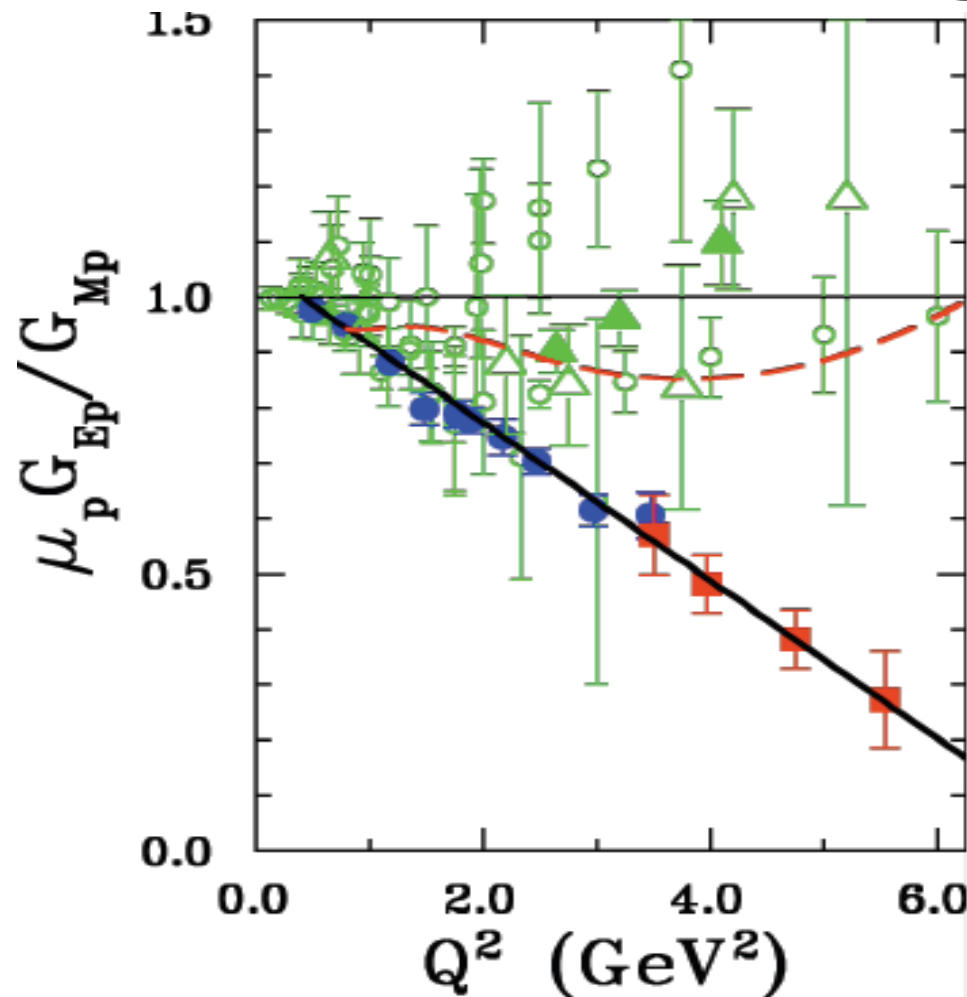
$$\pm hP_t = \mp h 2\sqrt{\tau(1+\tau)} G_E^p G_M^p \tan\left(\frac{\theta_e}{2}\right) / I_0$$

$$\pm hP_l = \pm h(E_e + E_{e'}) (G_M^p)^2 \sqrt{\tau(1+\tau)} \tan^2\left(\frac{\theta_e}{2}\right) / M / I_0$$

$$I_0 = \left\{ G_E^p(Q^2) \right\}^2 + \tau \left\{ G_M^p(Q^2) \right\}^2 \left[1 + 2(1+\tau) \tan^2\left(\frac{\theta_e}{2}\right) \right]$$

$$\frac{G_E^p}{G_M^p} = -\frac{P_t}{P_l} \frac{E_e + E_{e'}}{2M} \tan\left(\frac{\theta_e}{2}\right)$$

Hall A GEp results



- Almost linear drop of G_{Ep}/G_{Mp} with Q^2 above.
- May be an indication of the importance of quark Orbital Angular Momentum
- Recent **super-Rosenbluth measurement** confirmed old result: " $\mu G_{Ep}/G_{Mp}$ " = 1
- The difference is believed to be due to **Two-photon exchange** corrections.
- Several experiments to check this hypothesis.

- **E93-027** PRL 84, 1398 (2000)
Used both HRS in Hall A with FPP
- **E99-007** PRL 88, 092301 (2002)
used Pb-glass calorimeter for electron detection to match proton HRS acceptance

Important role of quark Orbital Angular Momentum ?

pQCD (Bjorken) scaling, taken with the assumption of Hadron Helicity Conservation predicts

Conservation predicts

$$F_1 \propto 1/Q^4 ; F_2 \propto 1/Q^6$$

$$\Rightarrow F_2/F_1 \propto 1/Q^2 \text{ (Brodsky \& Farrar)}$$

Data clearly do not follow this trend

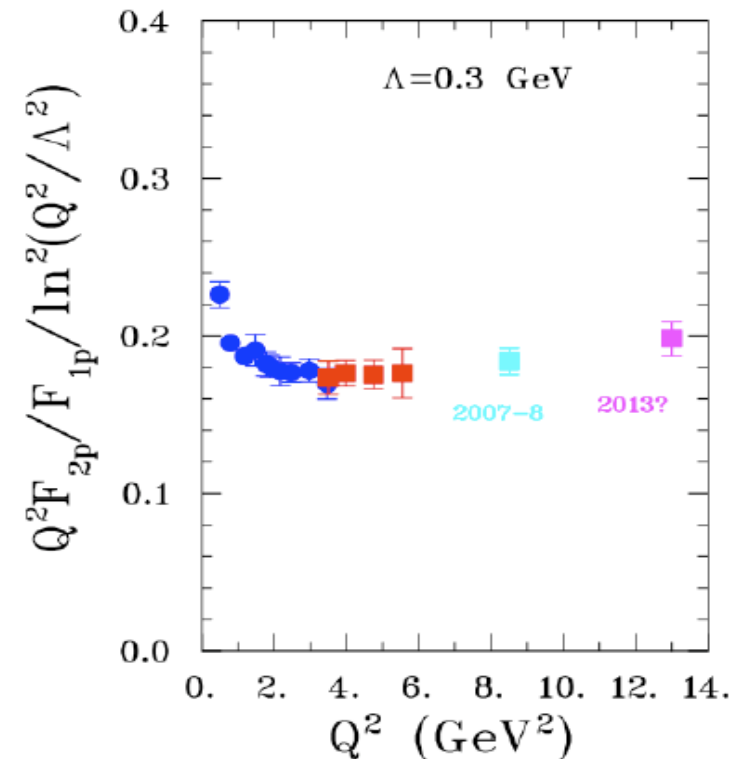
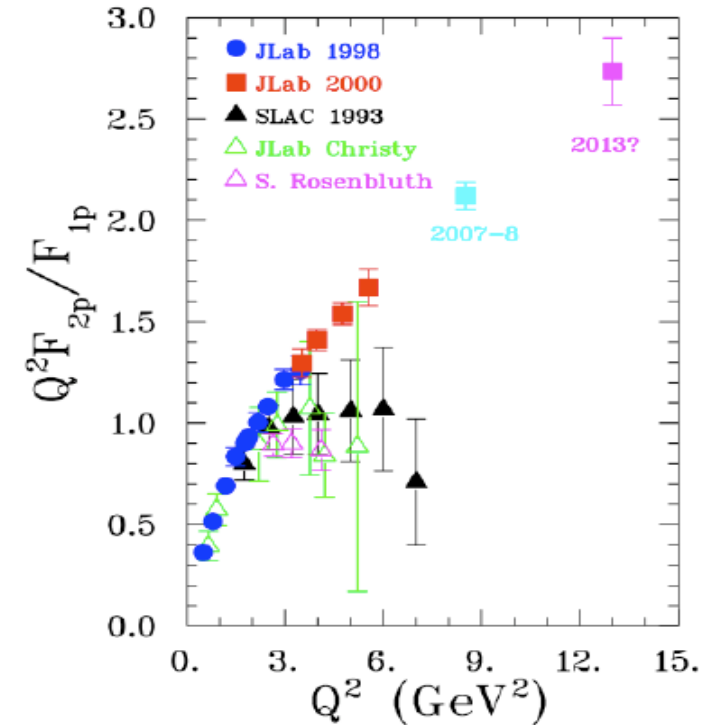
Schlumpf (1994), Miller (1996) and Ralston (2002) agree that by

- freeing the $p_T=0$ pQCD condition
- applying a (Melosh) transformation to a relativistic (light-front) system
- an orbital angular momentum component is introduced in the proton wf (giving up helicity conservation) and one obtains

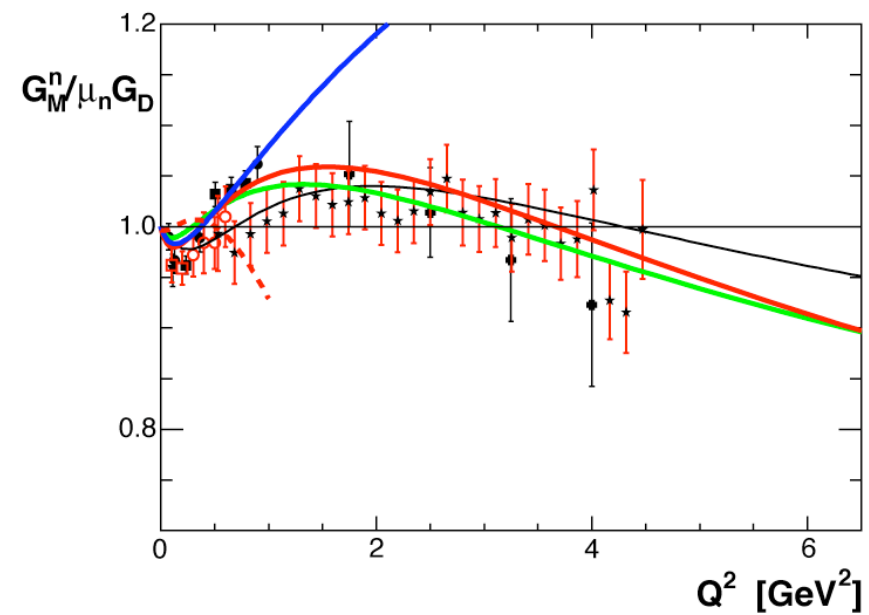
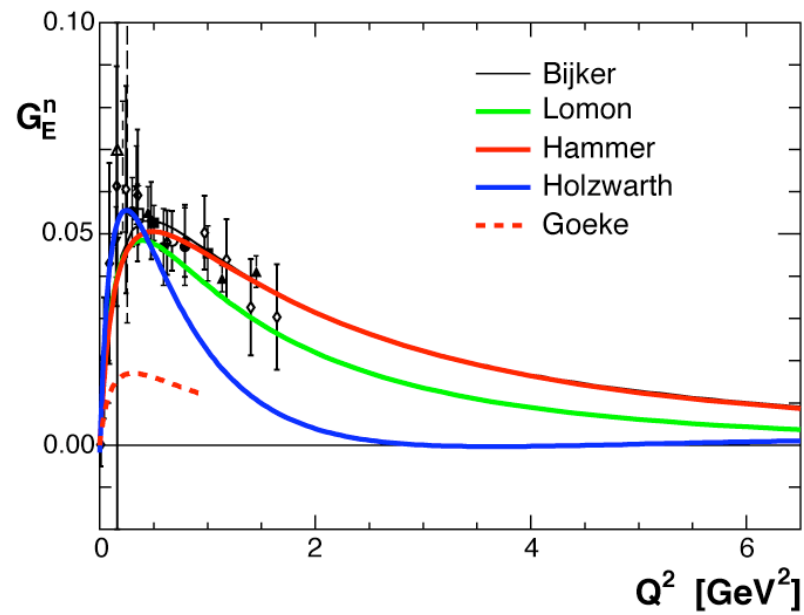
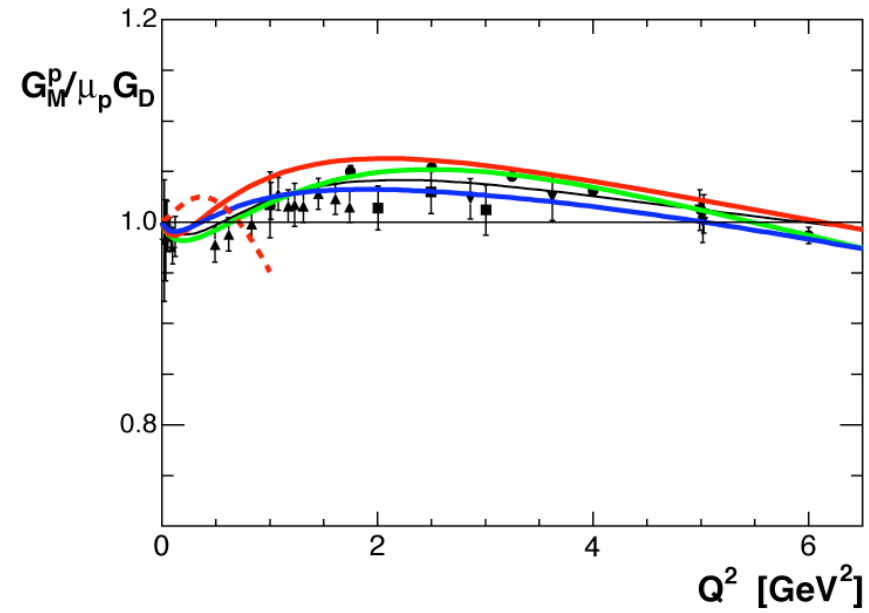
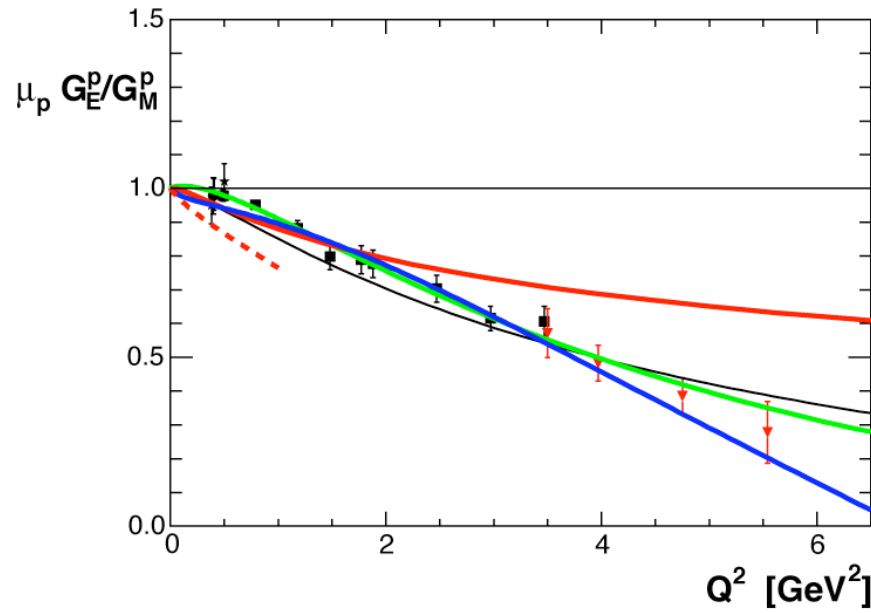
$$\Rightarrow F_2/F_1 \propto 1/Q$$

- and equivalently a linear drop off of G_E/G_M with Q^2
- Belitsky, Ji and Yuan refined the pQCD prediction by including quark OAM component $l_z=1$ (relax HHC) and get:

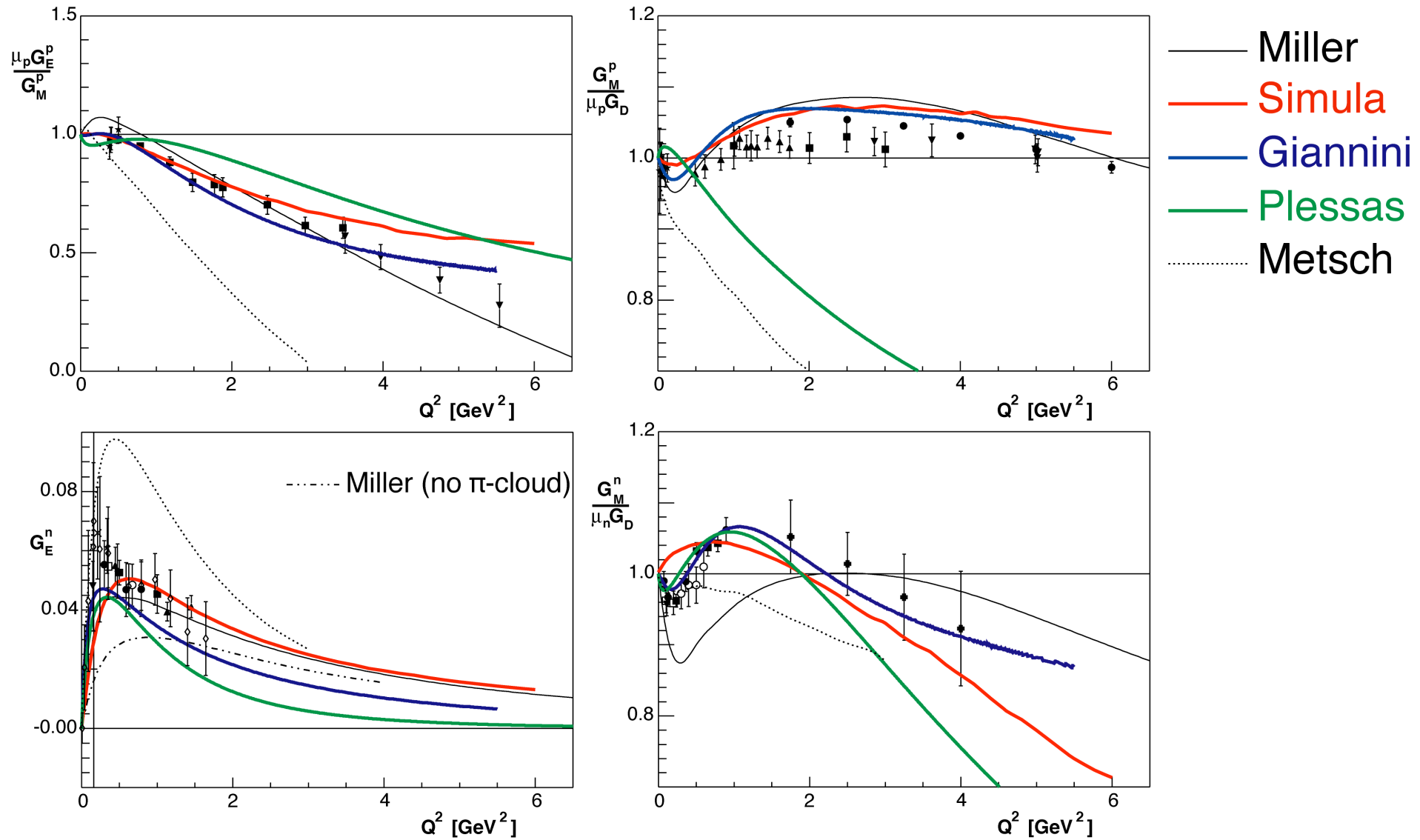
$$F_2/F_1 \propto \ln^2(Q^2/\Lambda^2)/Q^2$$



Vector-Meson Dominance models



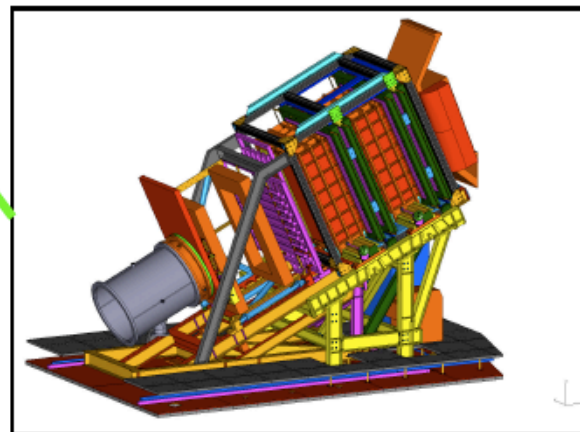
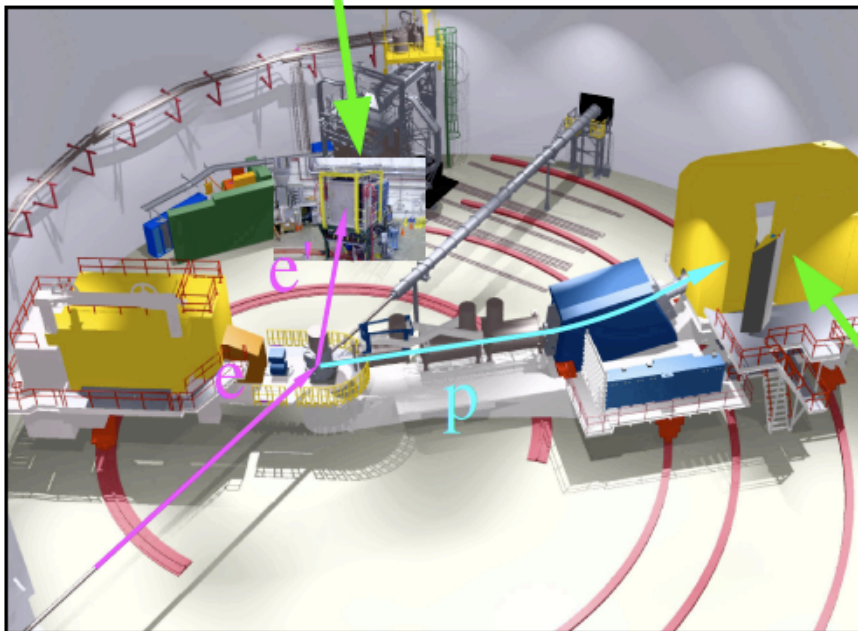
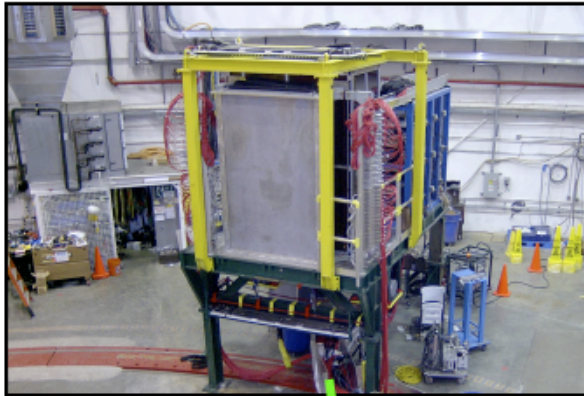
Relativistic Constituent Quark models



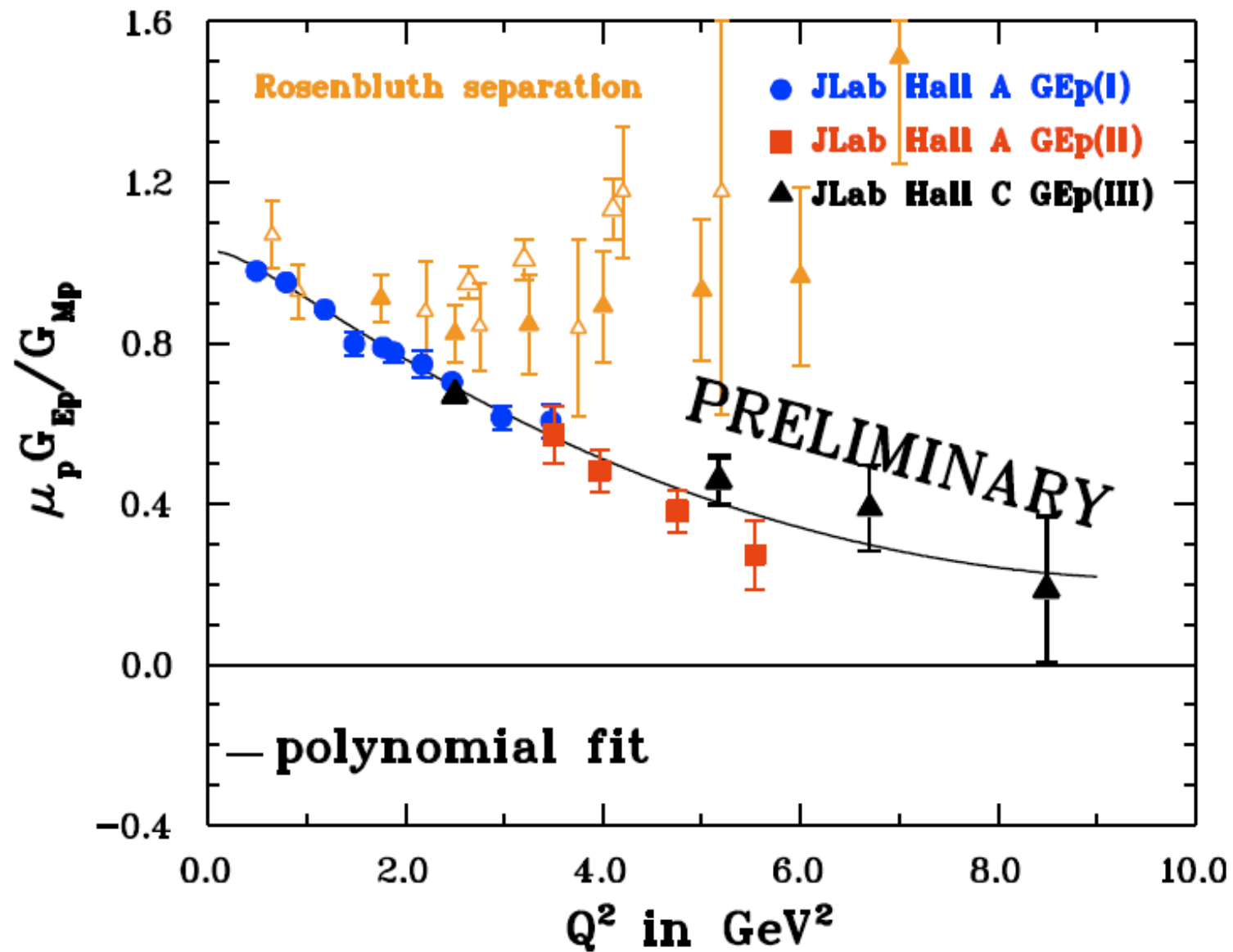
Need electric and magnetic form factor data for both neutron and proton up to high Q^2 with the highest possible precision in order to to discriminate between models and to understand underlying nucleon dynamics

Jefferson lab Hall C GEP-III experiment (E04-108)

- Continuation of Hall A Gep 1+2 experiments to higher Q^2 : 5.2, 6.8, and 8.5 GeV^2
- Elastic electron detected in large Calorimeter, BigCal
- Proton detected in HMS equipped with new FPP.
- The first high Q^2 polarization transfer Gep measurement outside Hall A.



Preliminary results of GEp(III)

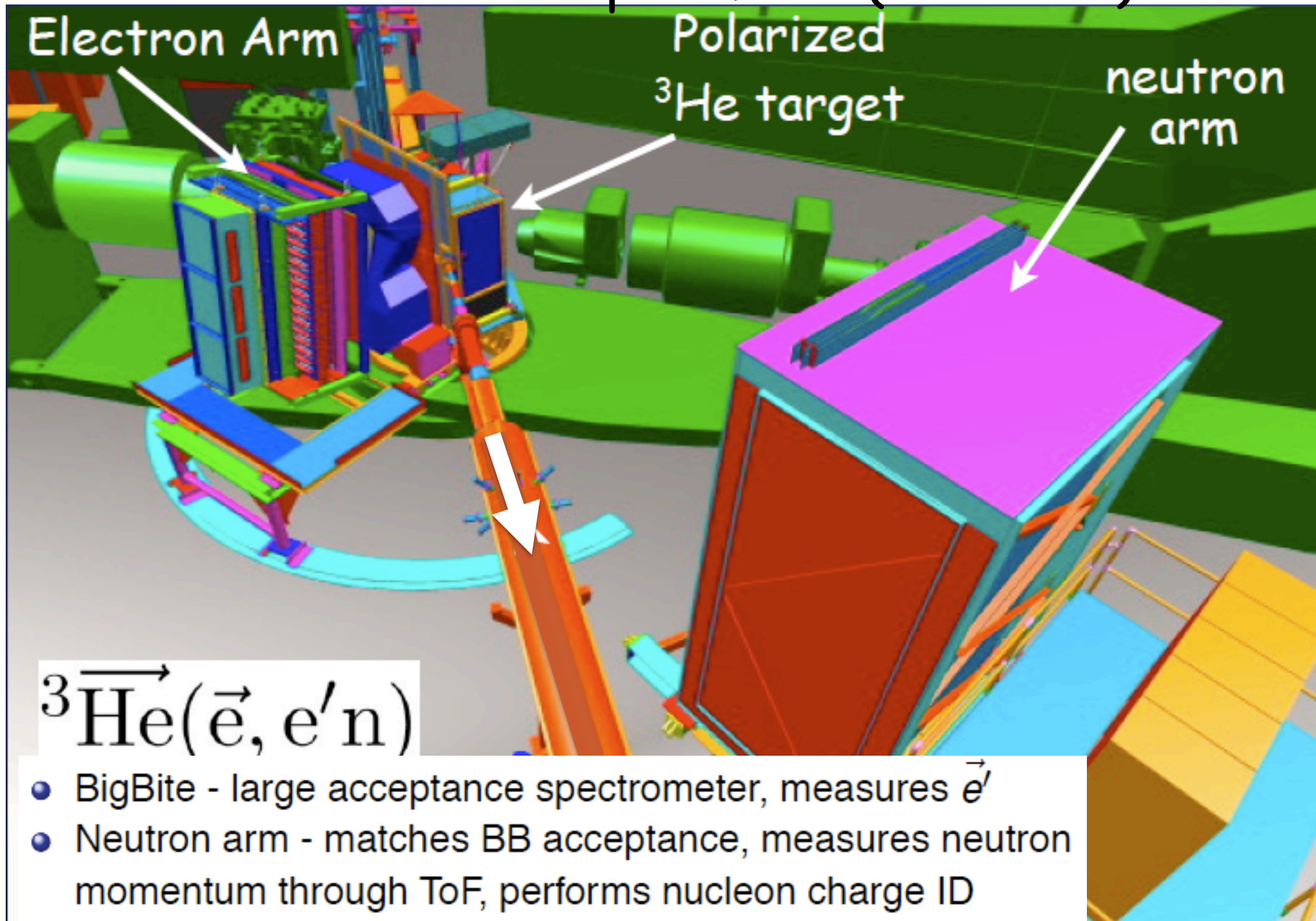


Double Polarization method to measure G_E

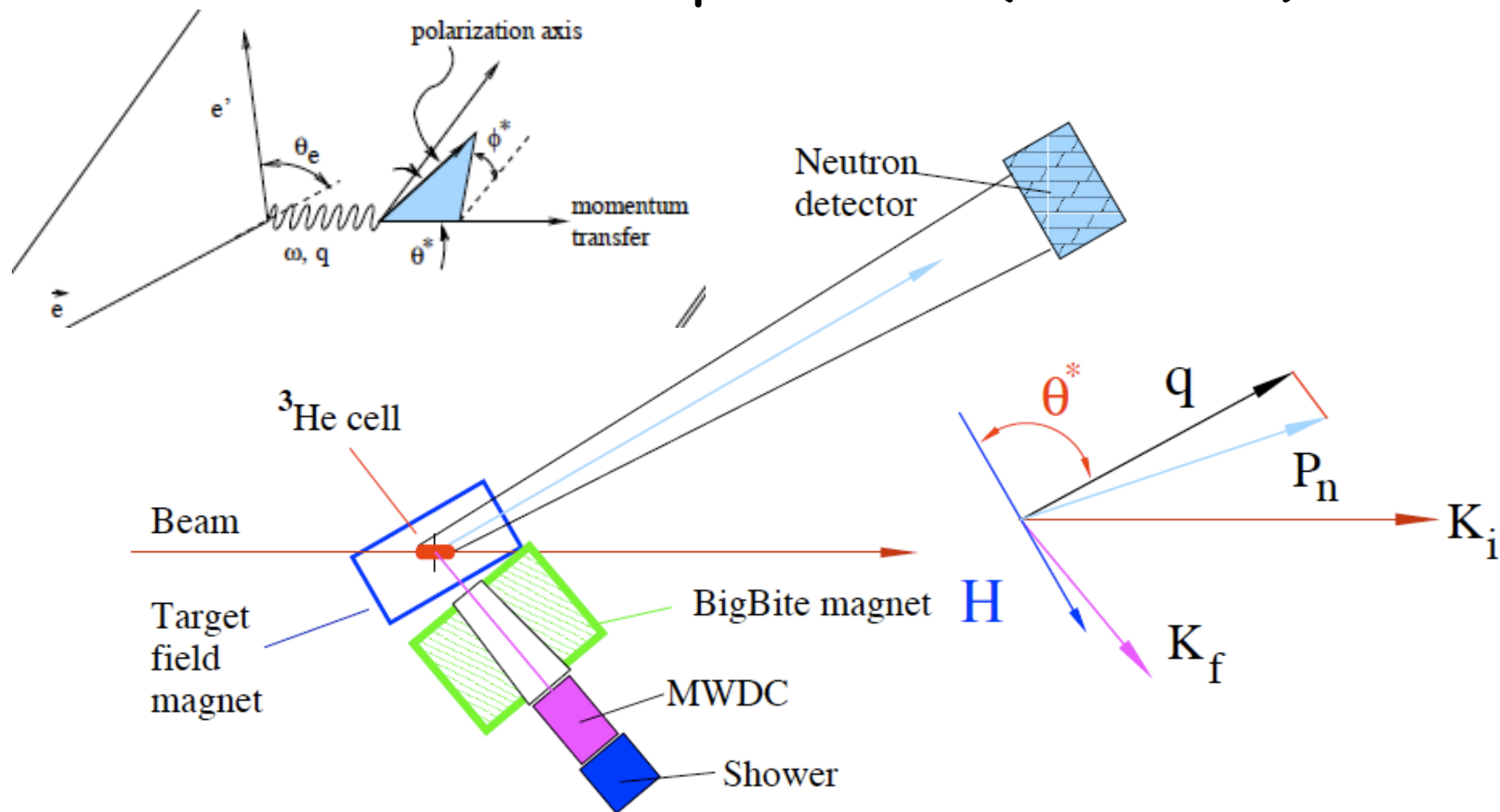
- Use $(e,e'n)$ reaction off deuterium with polarized target (like ND_3) or with recoil polarimetry.
 - Problems
 - Low deuteron polarization ($\sim 25\%$)
 - Low beam on target ($< 100 \text{ nA}$)
 - Low FOM for neutron polarimeters.
- Use $(e,e'n)$ off polarized ^3He target
 - Recent advances make this method very attractive
 - $> 50\%$ target polarization with over $12 \mu\text{A}$ of beam
 - Above $Q^2 \sim 1\text{-}2 \text{ GeV}^2$ Glauber method provides sufficiently accurate nuclear corrections.

Polarized ^3He target combined with a large acceptance electron spectrometer and a large neutron detector allows G_E to be measured to high Q^2 .

Hall A GEn experiment (E02-103)



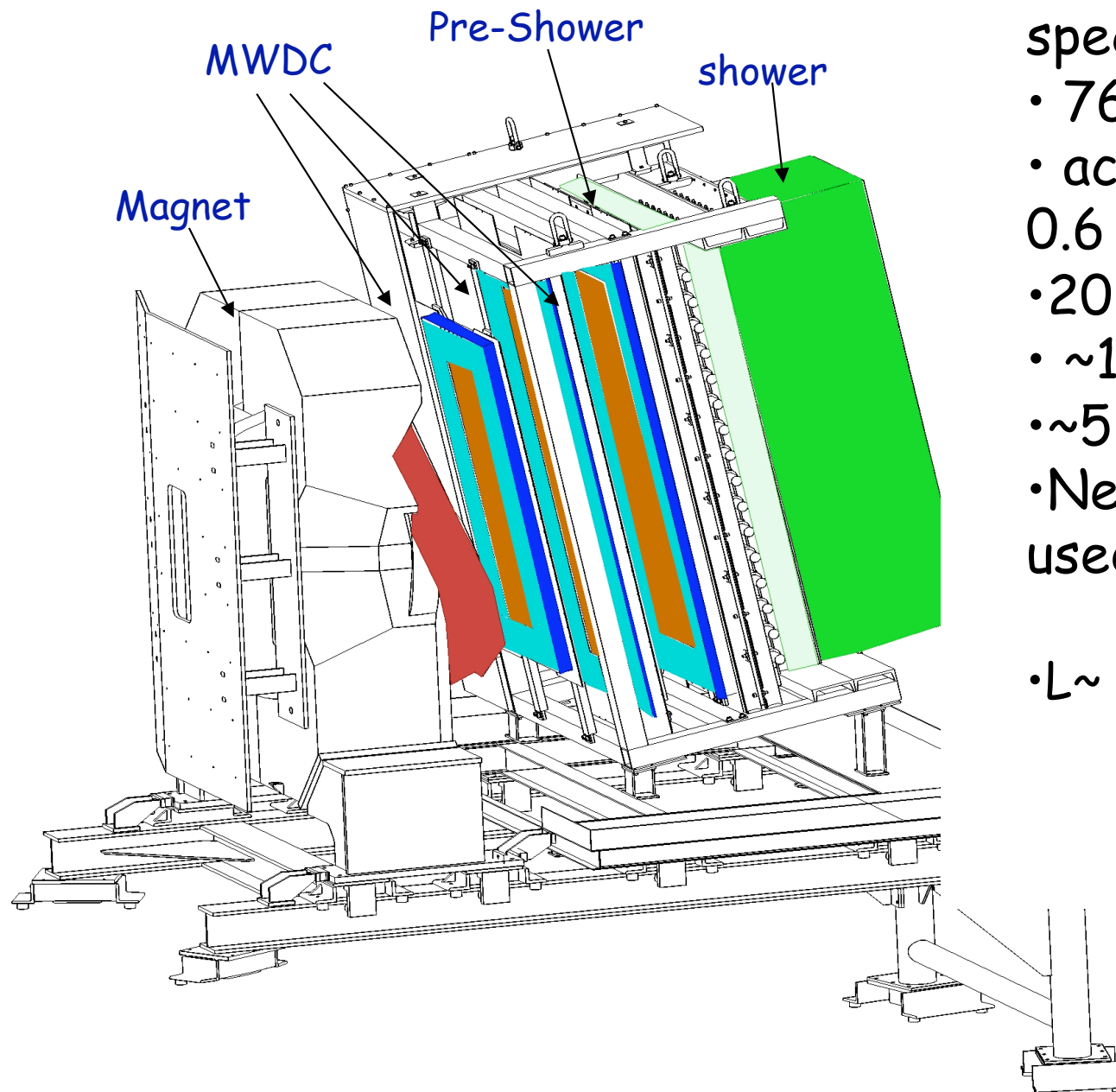
Hall A GEn experiment (E02-103)



$$A_{phys} = A_{\perp} + A_{\parallel} = \frac{a \cdot (G_E/G_M) \sin \theta^* \cos \phi^*}{(G_E/G_M)^2 + c} + \frac{b \cdot \cos \theta^*}{(G_E/G_M)^2 + c}$$

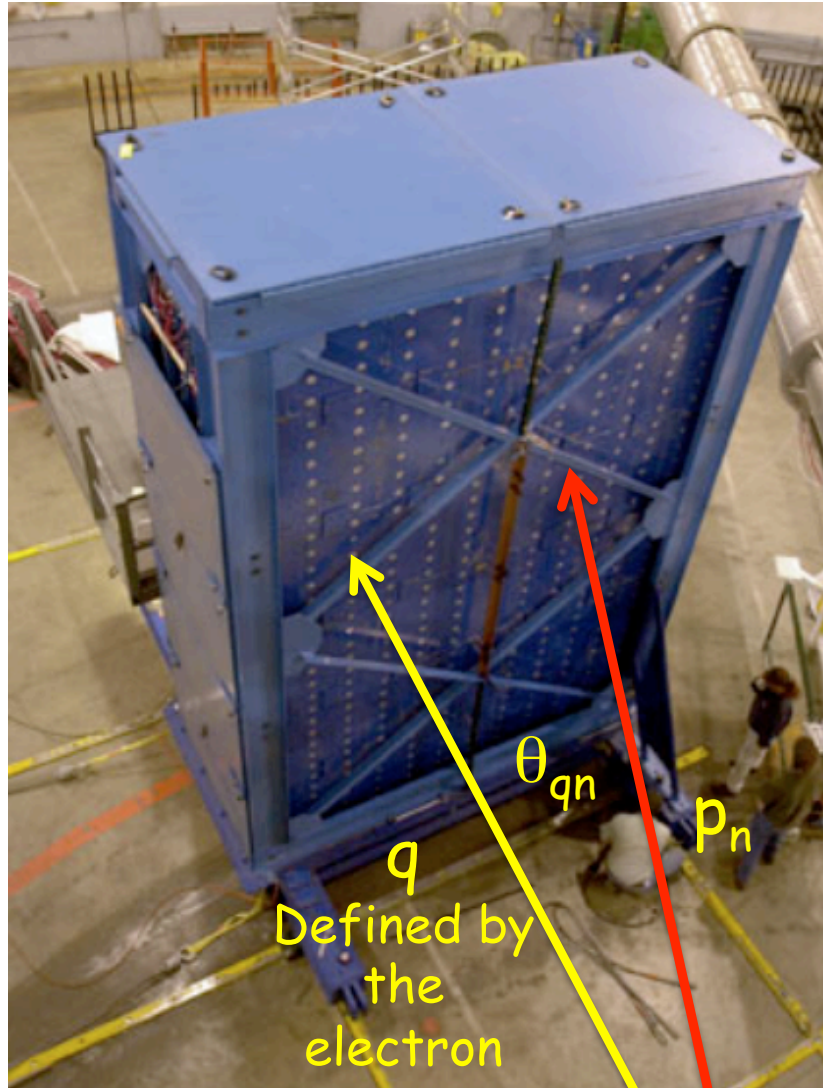
Here a , b and c are solely functions of kinematic factors (and not θ^* or ϕ^*)

Bigbite Spectrometer

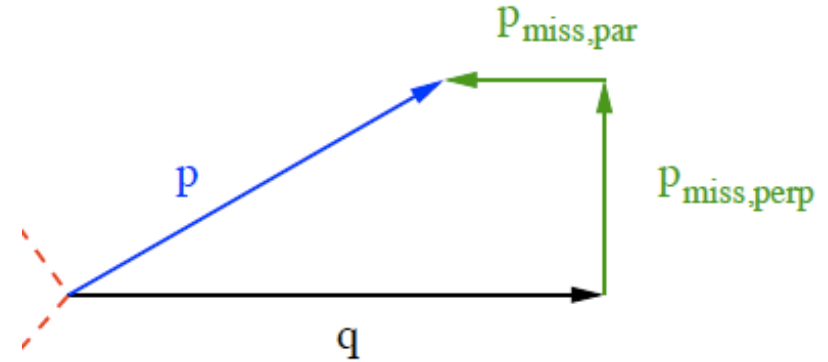


- Non-focusing large angular and momentum acceptance spectrometer
- 76 msr over 40 cm of target.
- accepting electrons between 0.6 ~ 1.6 GeV
- 20 MHz/plane on MWDC.
- ~1% momentum resolution
- ~5 mm y_{tg} resolution
- New detector package, first used for Gen
- $L \sim 4.5 \times 10^{36} \text{ cm}^{-2}\text{s}^{-1}$

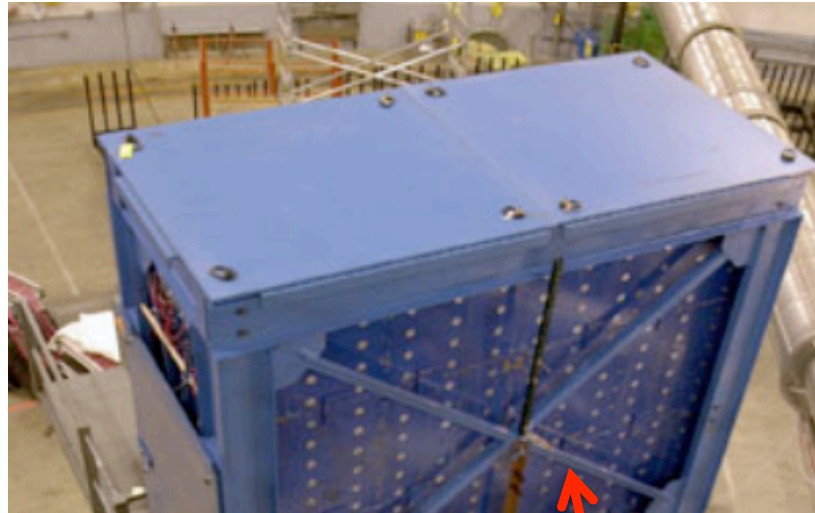
Neutron Detector - BigHAND



- Match BigBite solid angle in QE kinematics.
- Flight-path ~ 10 m
- 1.6×5 m² active area.
- 7 layers of scintillator bars + 2 layers of veto bars.
- time resolution ~ 0.3 ns.
- θ_{qh} is very effective in suppressing background:

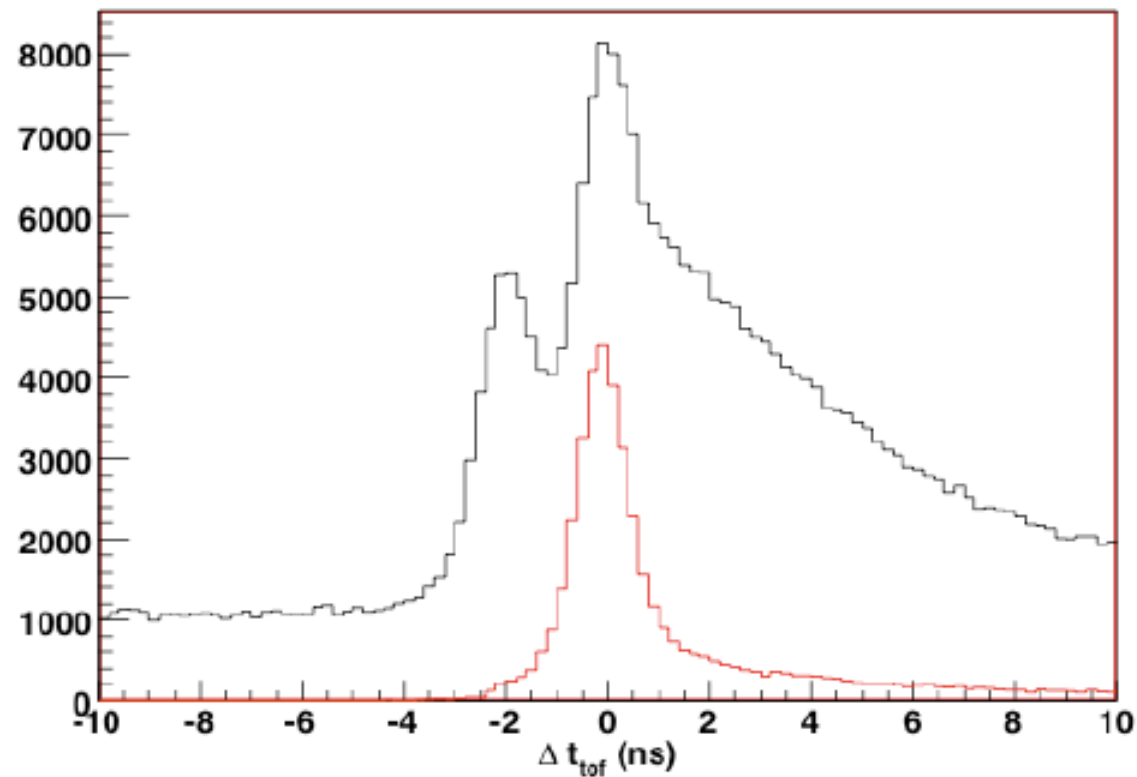


Neutron Detector - BigHAND

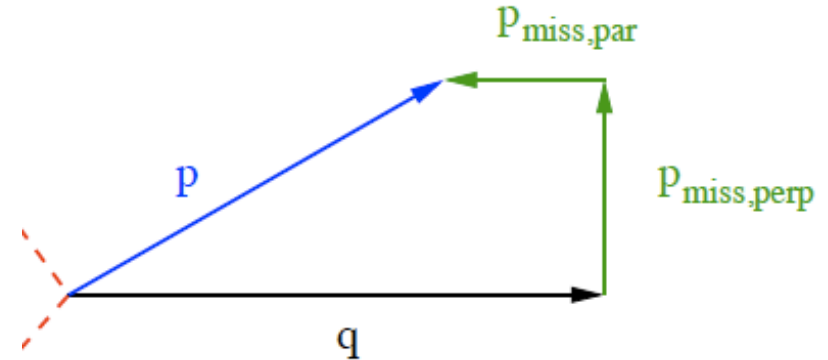


- Match BigBite solid angle in QE kinematics.
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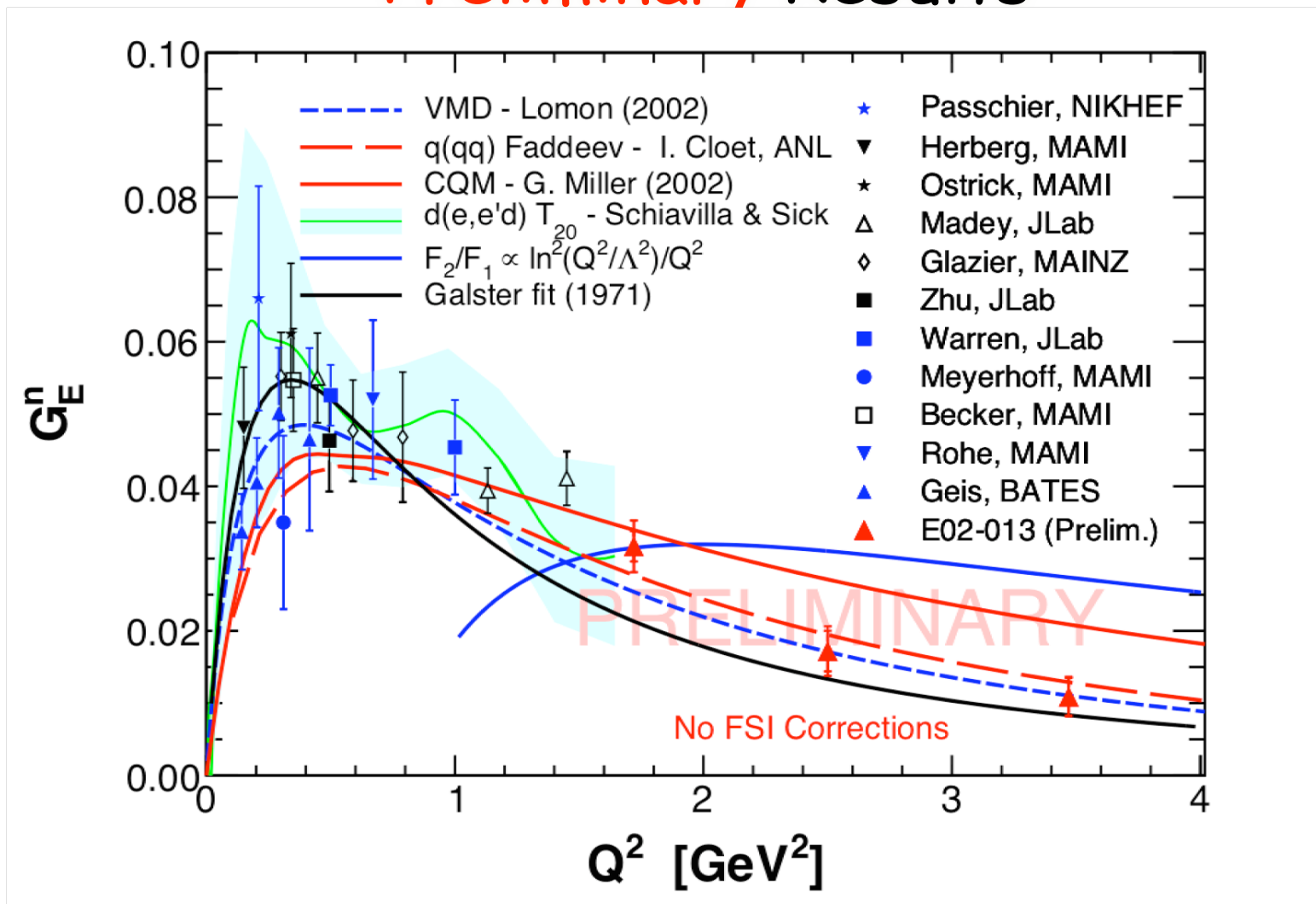
$\Delta t_{\text{tof}}, Q^2 = 3.5 \text{ GeV}^2$



very effective in
reducing background:



Preliminary Results



Still very Preliminary:

still to be done:

- Neutron detector simulation need to be finalized: will select final cuts based on this
- Nuclear corrections - neutron polarization and FSI

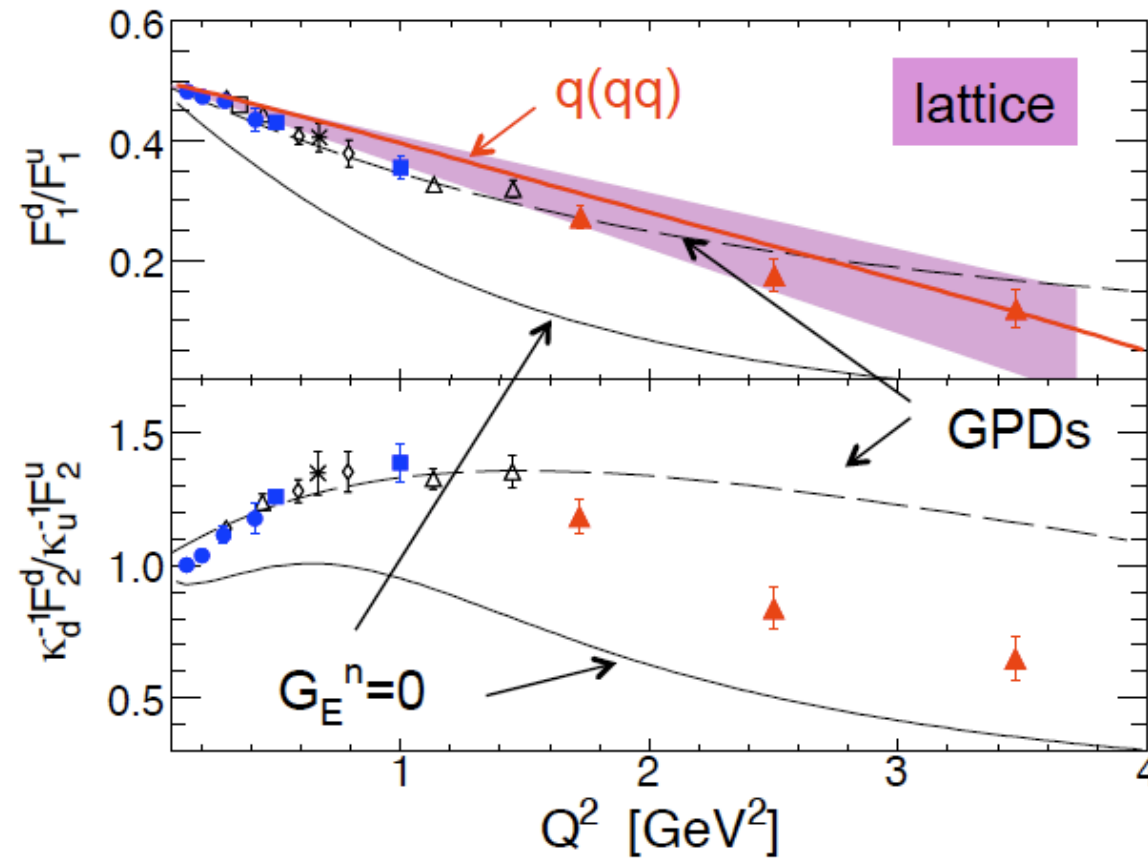
$F_{1(2)}^d/F_{1(2)}^u$ with proton and neutron FFs

$$F_1 = \frac{G_E + \tau G_M}{1 + \tau}$$

$$F_2 = -\frac{G_E - G_M}{1 + \tau}$$

$$F_1^u = 2F_{1p} + F_{1n}$$

$$F_1^d = 2F_{1n} + F_{1p}$$



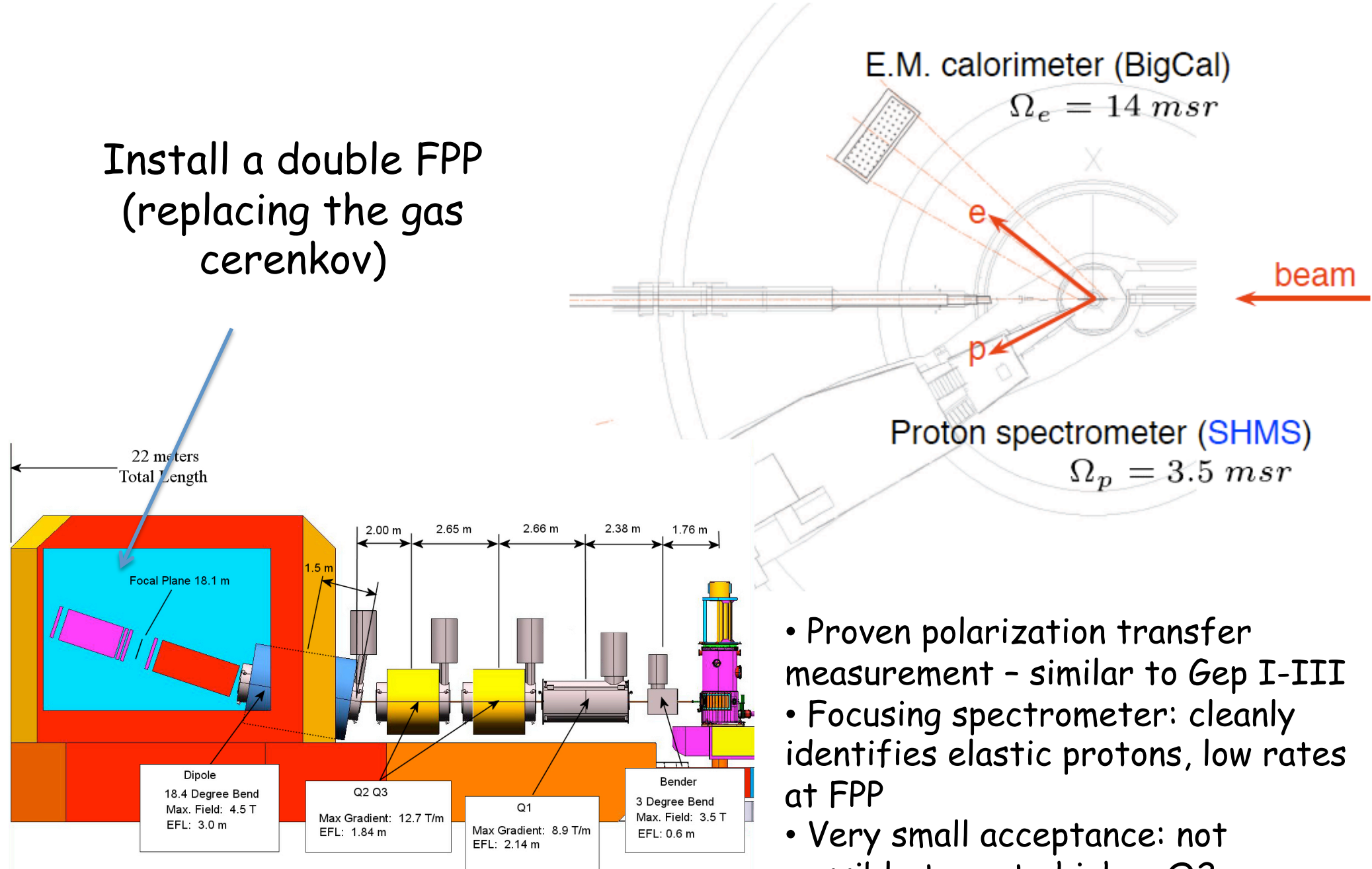
still to be done:

- Neutron detector simulation need to be finalized: will select final cuts based on this
- Nuclear corrections - neutron polarization and FSI

High Q^2 form factor measurements
with CEBAF 12 GeV beam

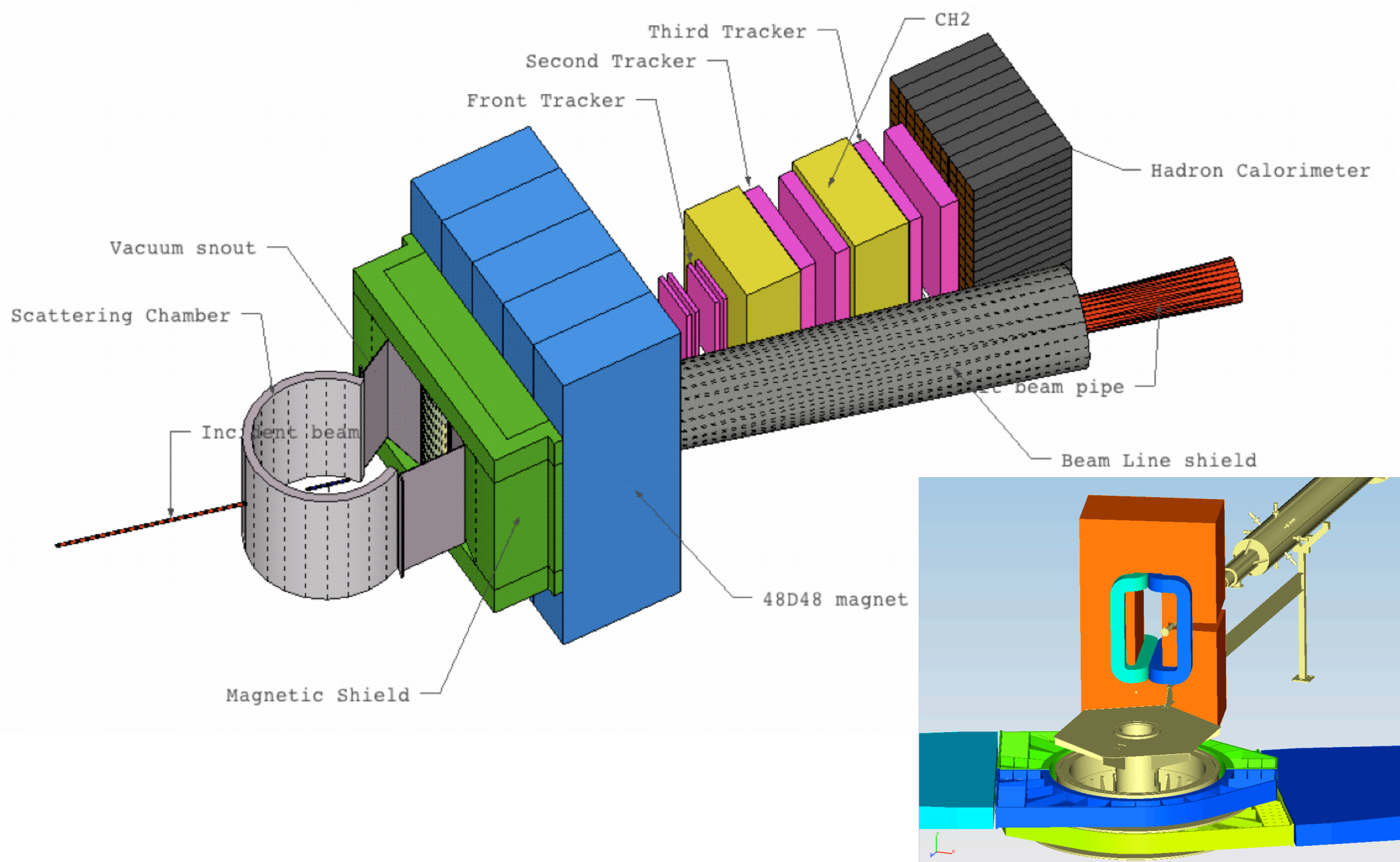
GEP/GMP up to Q2 = 13 GeV2 SHMS - Gep-IV

Install a double FPP
(replacing the gas
cerenkov)

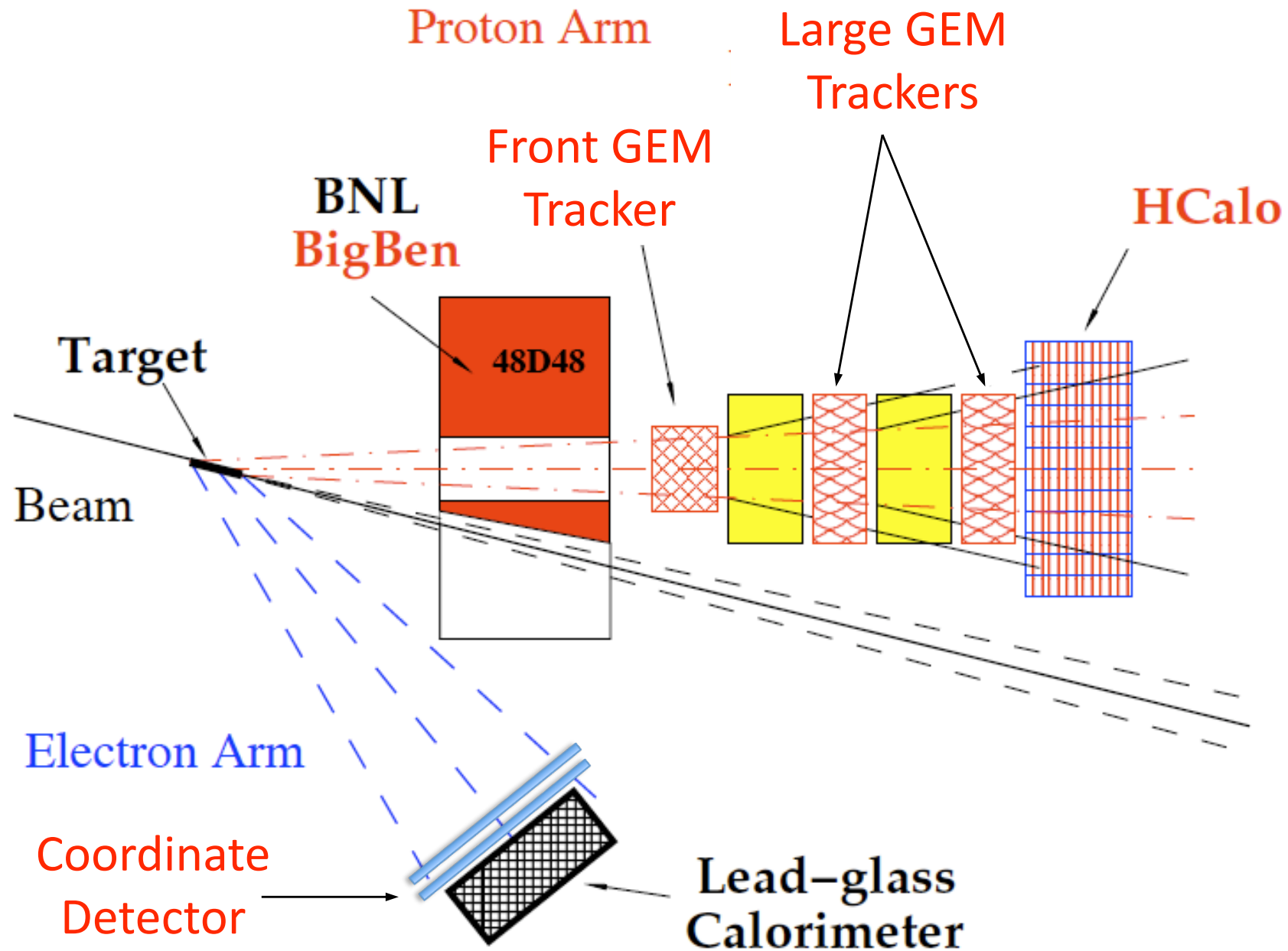


- Proven polarization transfer measurement - similar to Gep I-III
- Focusing spectrometer: cleanly identifies elastic protons, low rates at FPP
- Very small acceptance: not possible to go to higher Q2

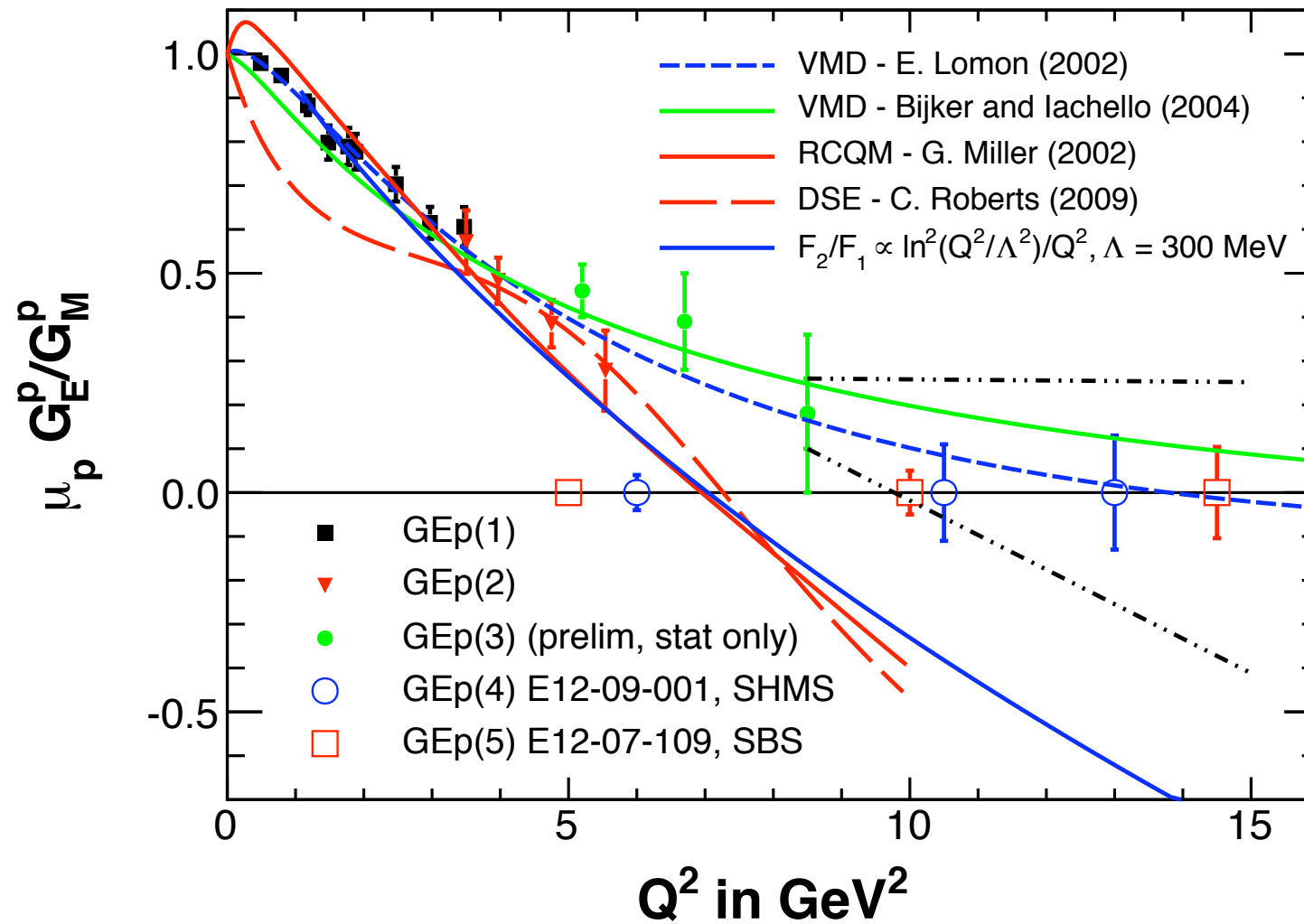
GEp, GEn and GMp measurements with Super-Bigbite spectrometer in Hall A



GEP/GMP up to $Q^2 = 15 \text{ GeV}^2$ with SBS

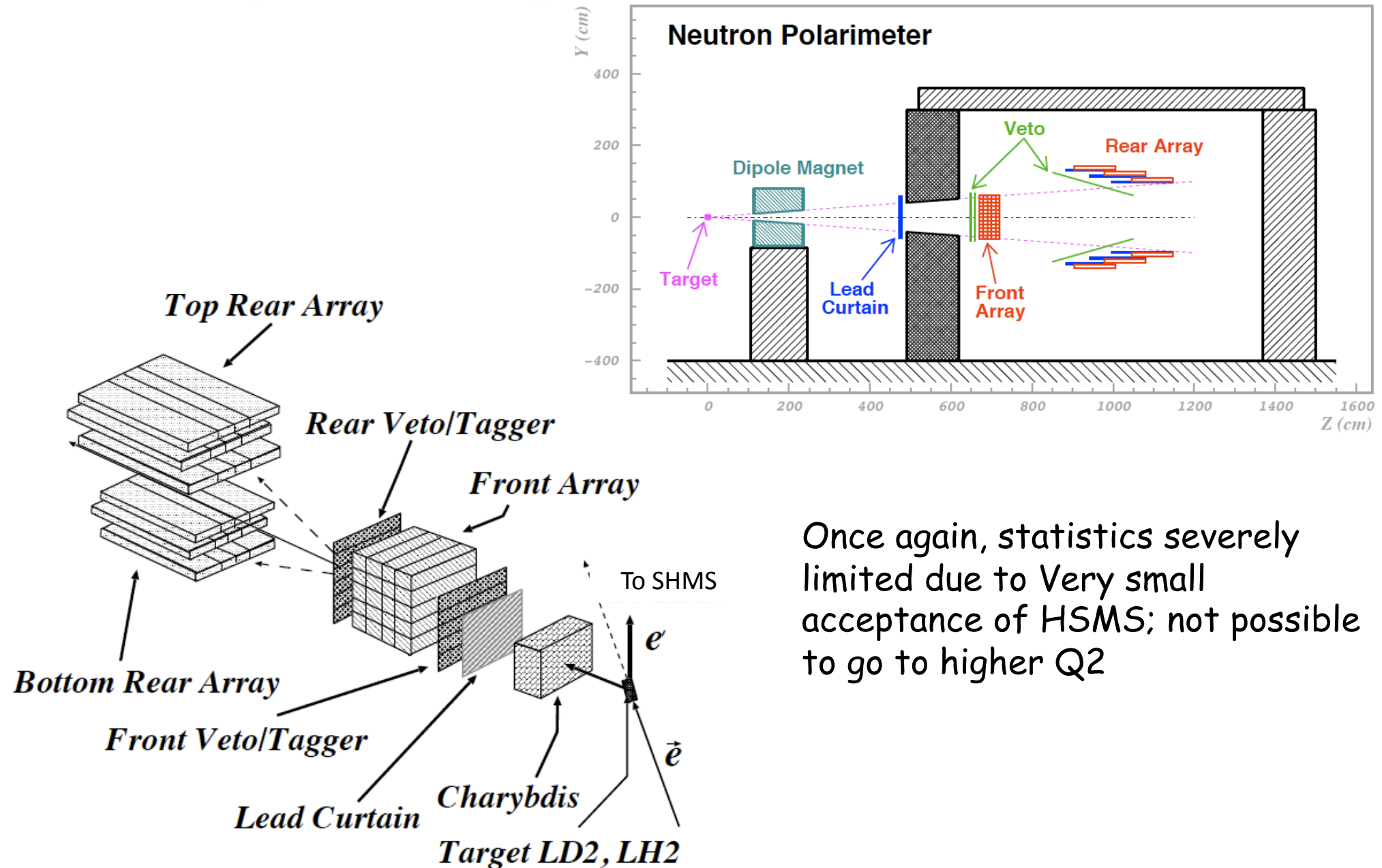


G_E^p/G_M^p up to $Q^2 = 15 \text{ GeV}^2$ with SBS



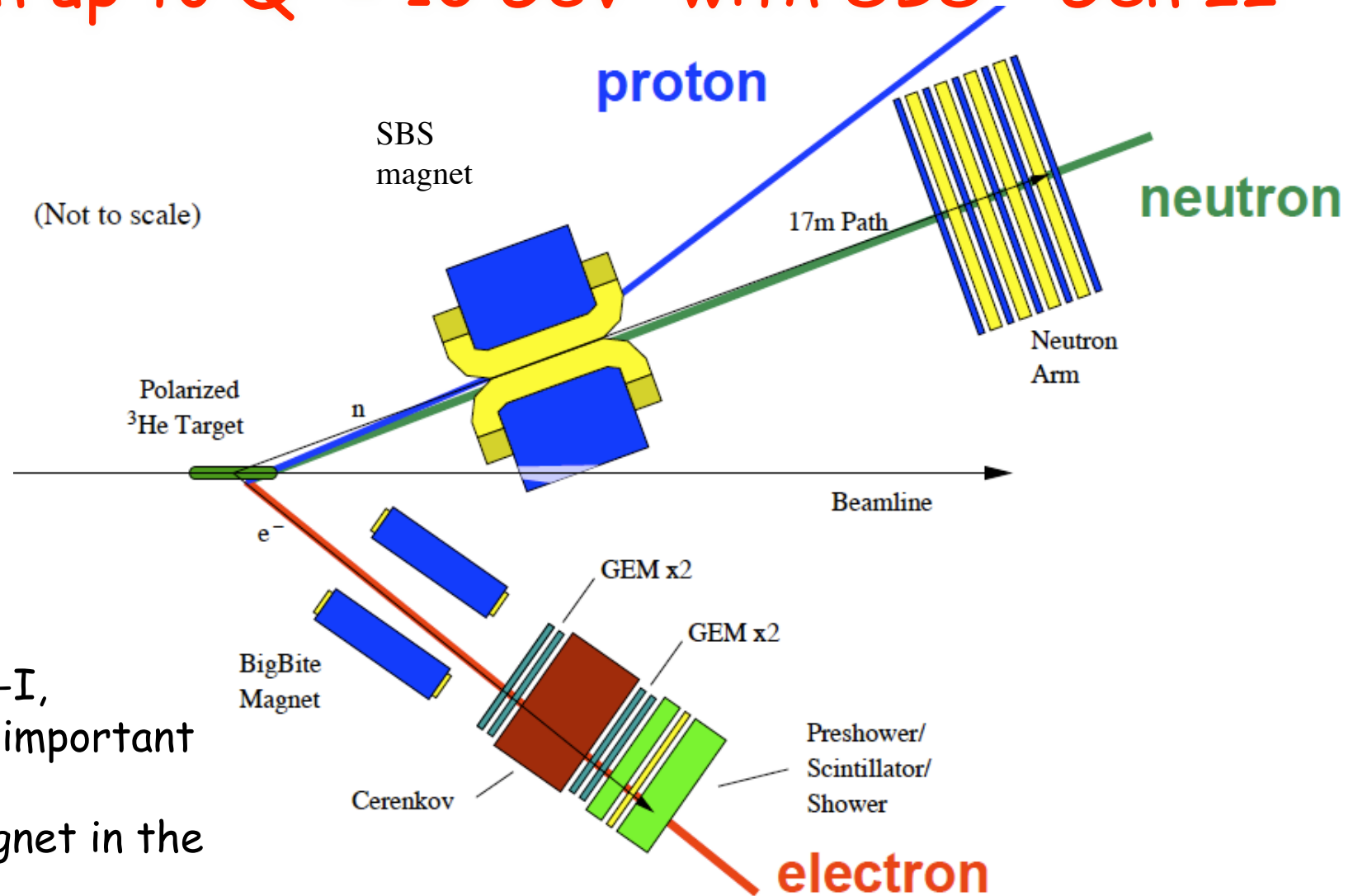
High Q^2 GEn Measurements with 12 GeV beam

E12-09-006: GEn up to $Q^2 = 7 \text{ GeV}^2$ using recoil polarimetry (similar to E93-038)



Once again, statistics severely limited due to Very small acceptance of HSMS; not possible to go to higher Q^2

GEn/GMn up to $Q^2 = 10 \text{ GeV}^2$ with SBS - Gen II



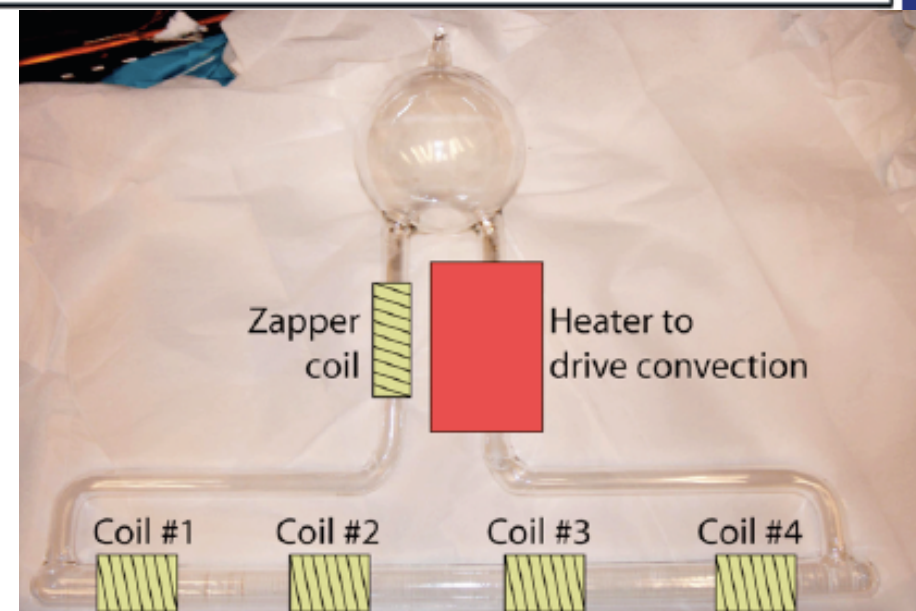
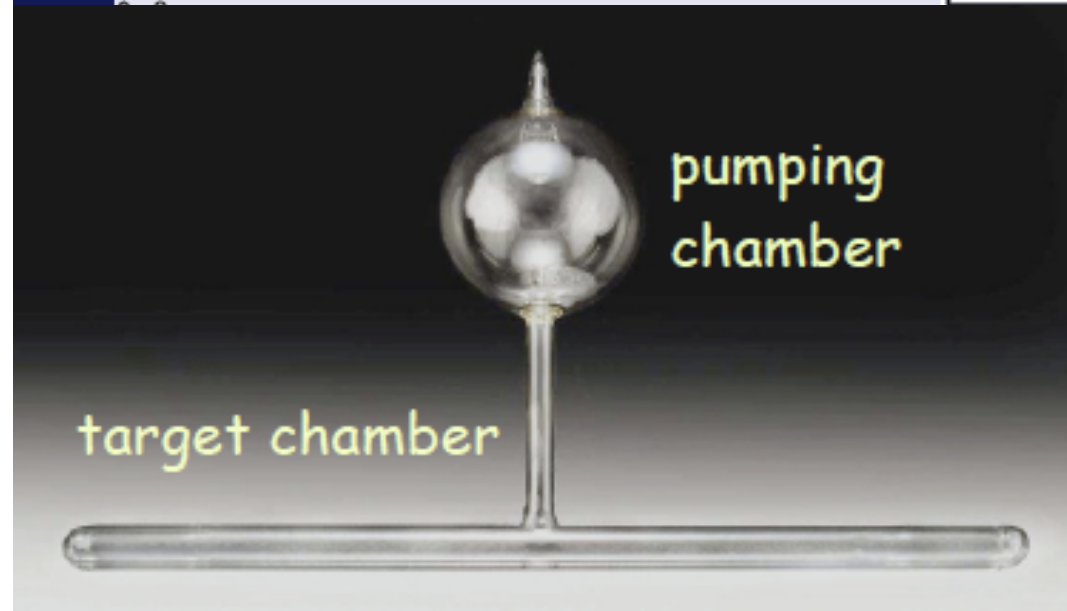
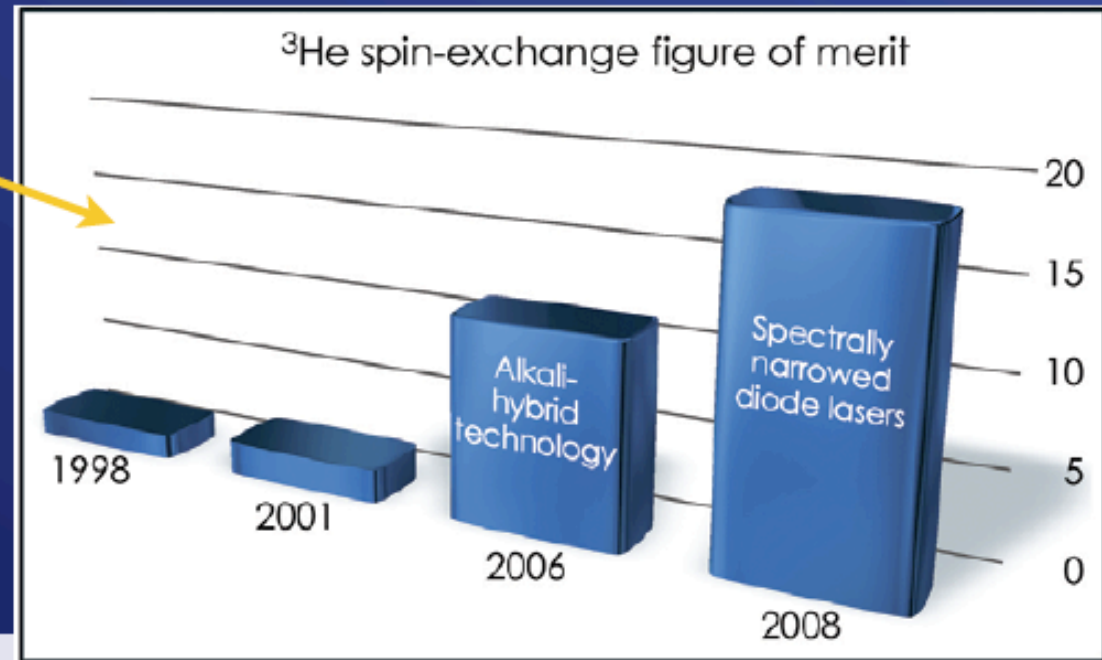
Similar to GEn-I,
But with some important
changes:

- Sweeping magnet in the Neutron arm
- Very high rate capable GEM-based Electron tracker + Cerenkov
- High-Luminosity pol. ^3He target: projected 60% polarization with $60 \mu\text{A}$ beam

Spin-exchange target will take advantage of huge improvements in F.O.M.

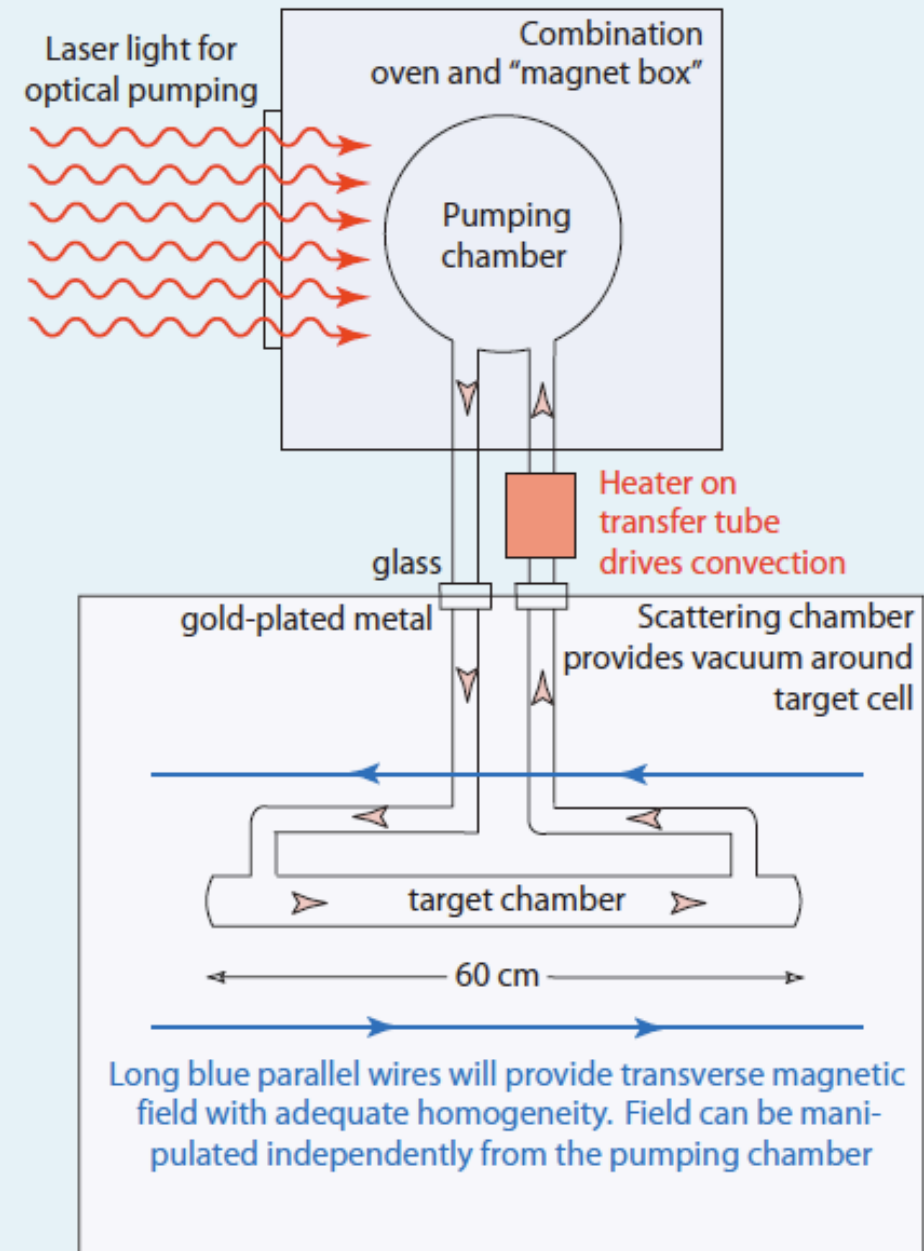
Target evolution at JLab

Shown is the F.O.M. defined as spins polarized per second weighted by polarization²

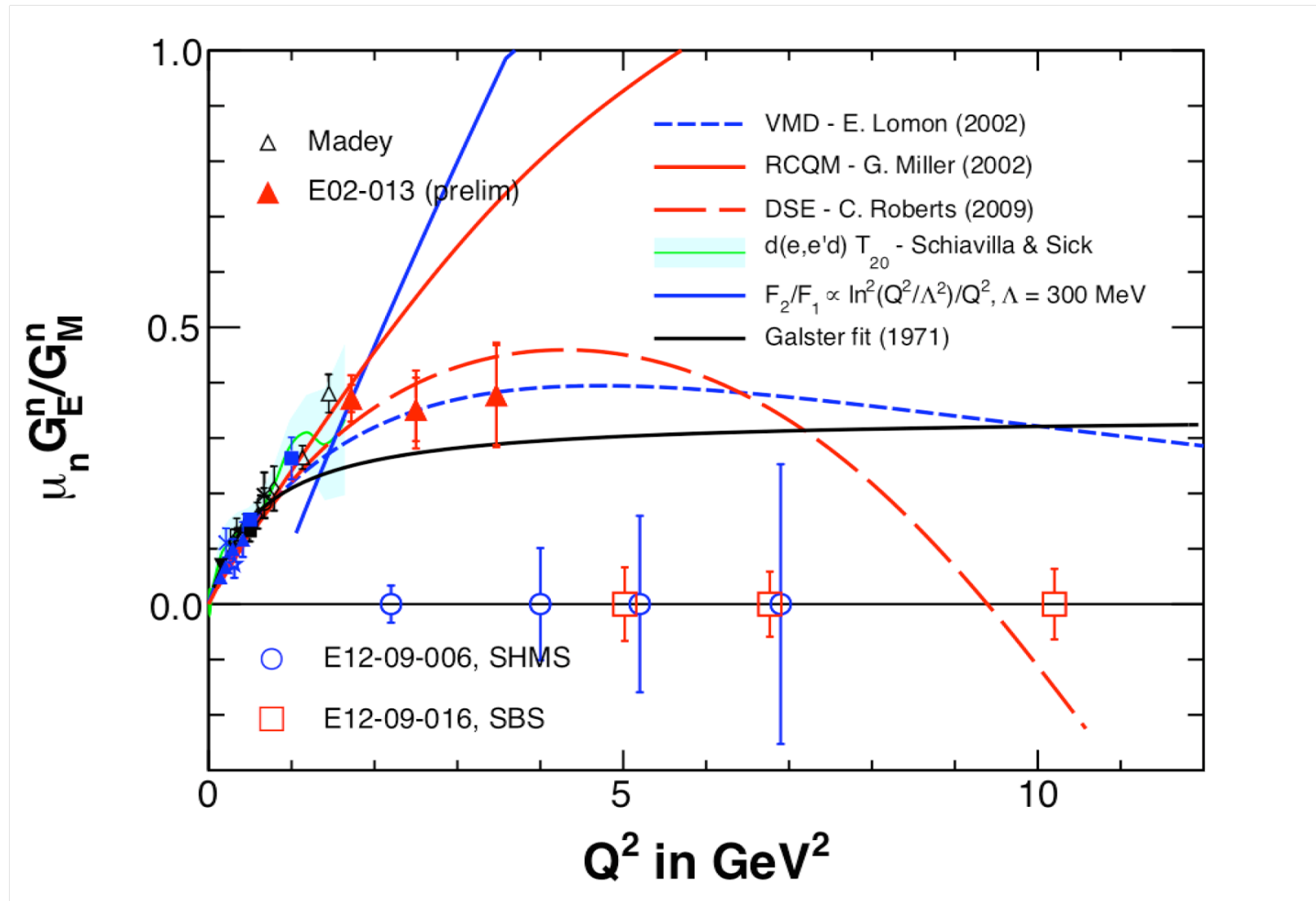


Very-high-luminosity polarized ^3He target

- Large pumping chamber provides ample reservoir of polarized spins to replenish the effects of intense electron beam.
- Convection-driven gas flow insures mixing times of minutes or less.
- Metal target cell (gold coated) ensures the target can physically tolerate the beam.
- Keeping the target cell in vacuum ensures the detectors see a manageable overall luminosity.
- Separating the pumping chamber and the target chamber by arbitrary distances greatly simplifies the magnetic field.

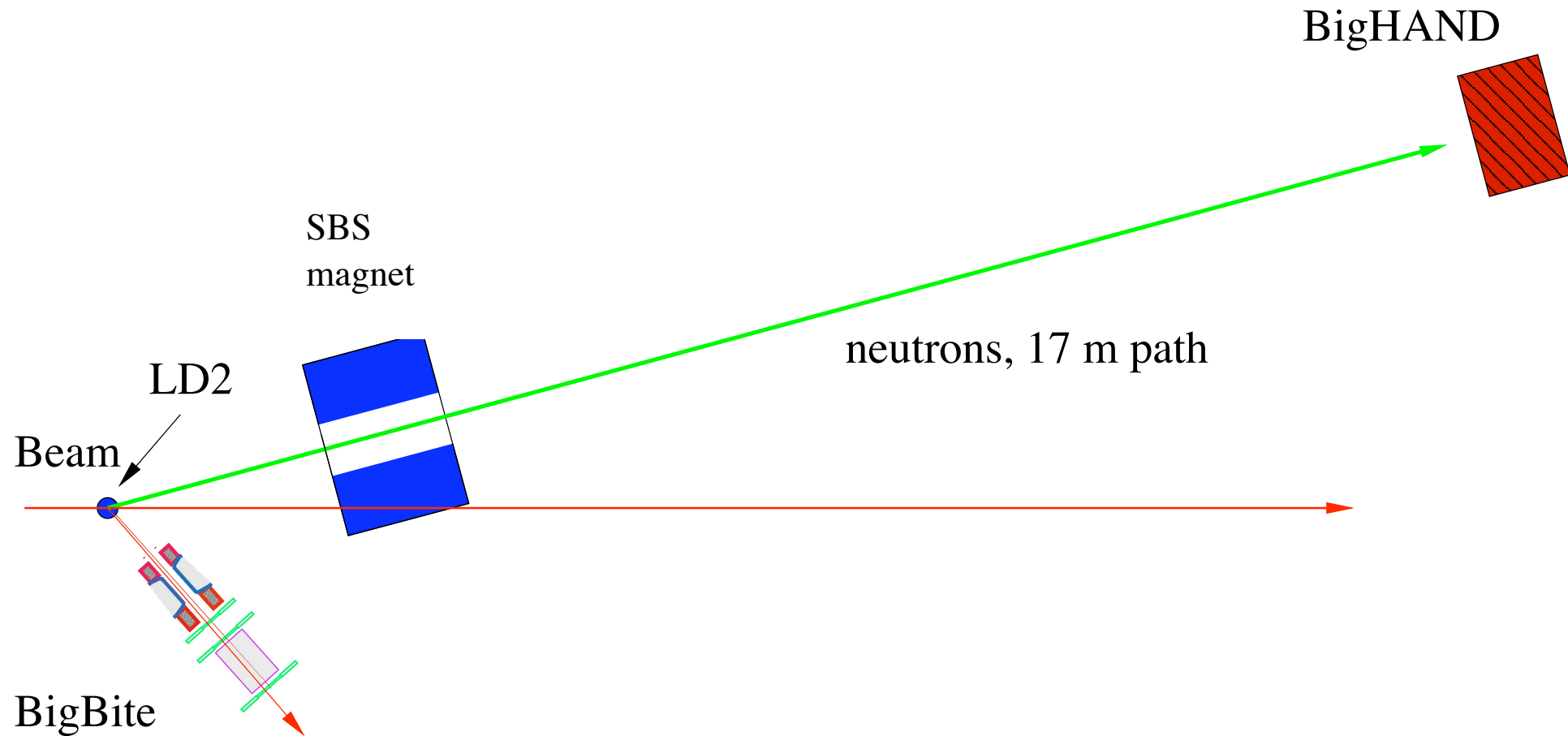


GEn/GMn up to $Q^2 = 10 \text{ GeV}^2$ with SBS - Gen II



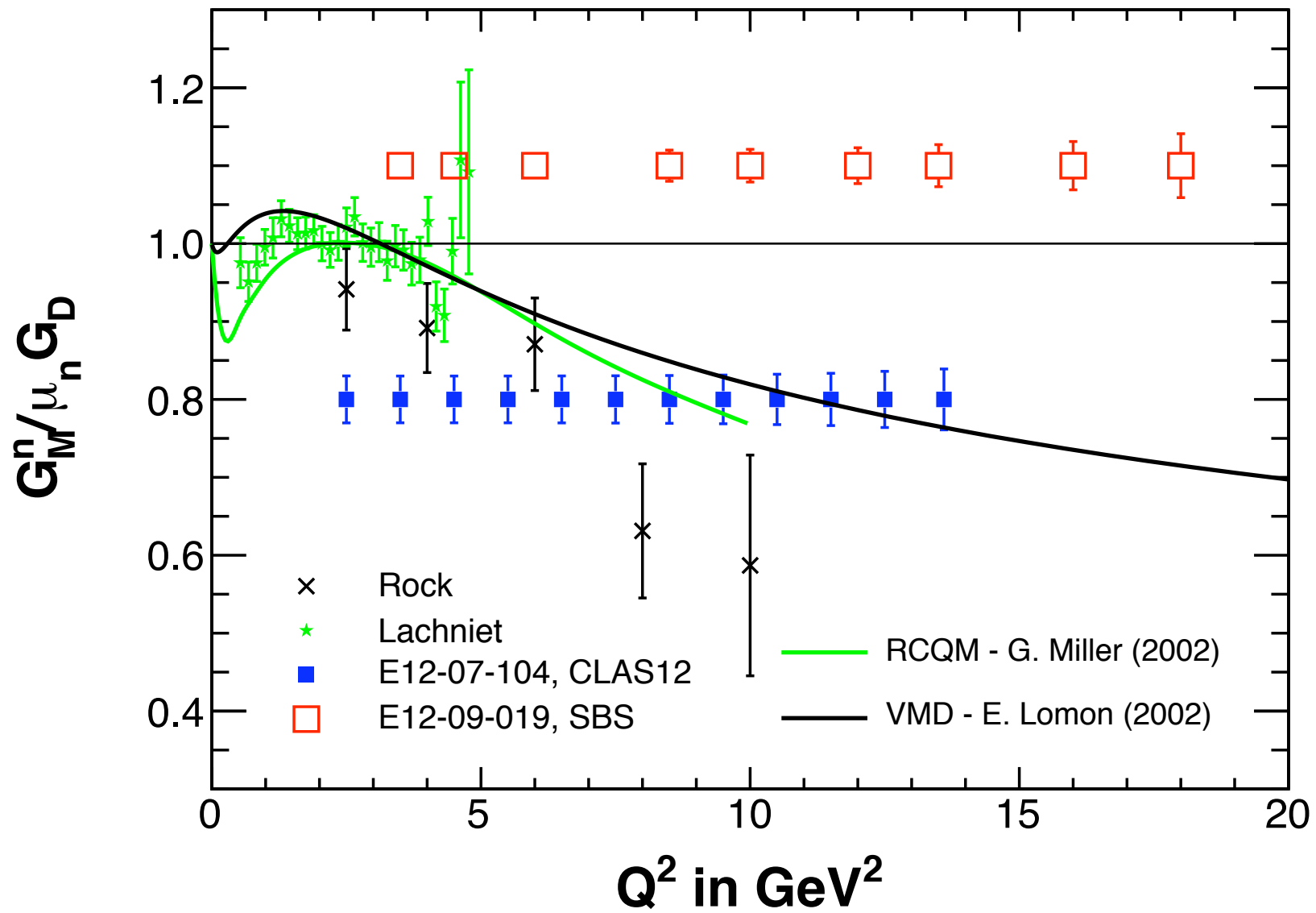
The SBS Gen experiment will provide clear discrimination between models (all of which seem to do an OK job fitting to high Q^2 GEp data)

GMn up to $Q^2 = 18 \text{ GeV}^2$ with SBS - Gen II



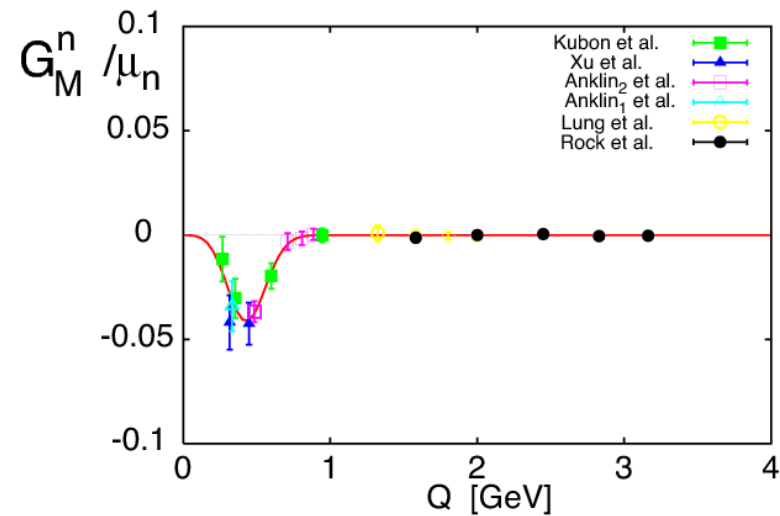
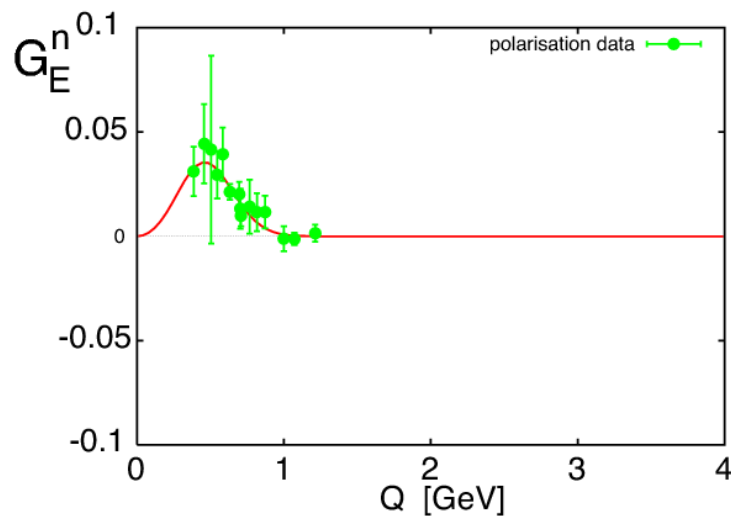
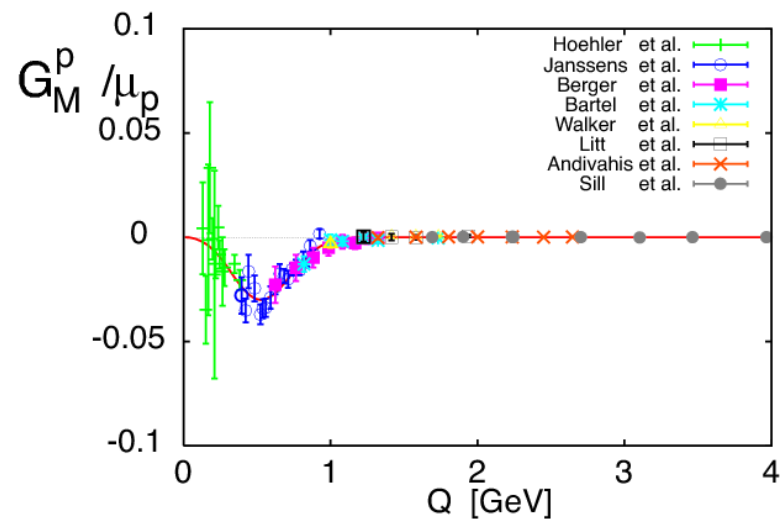
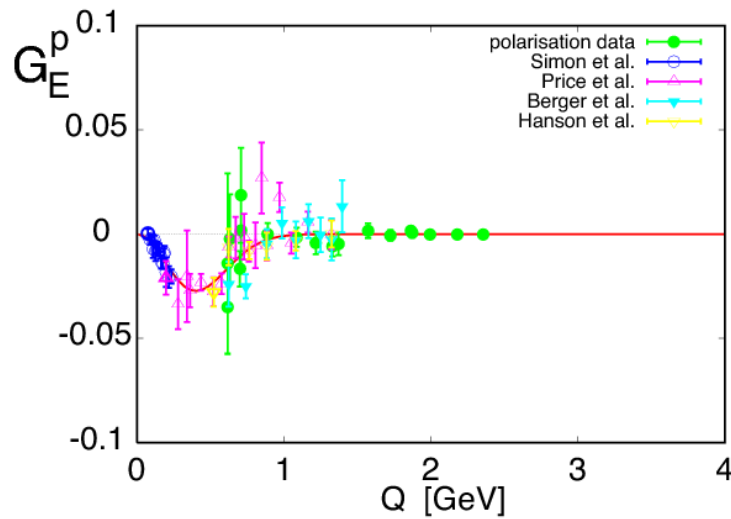
- Use the ratio method:
 $D(ee'n)/D(ee'p)$
- Use SBS magnet to kick the QE protons up; so they can be clearly separated from QE neutrons
- Both QE protons and neutrons are detected; for a given q -vector, the expected location of neutron is different from that of the proton
- Also has a veto-plane to assist with n/p identification

GMn up to $Q^2 = 18 \text{ GeV}^2$ with SBS - Gen II



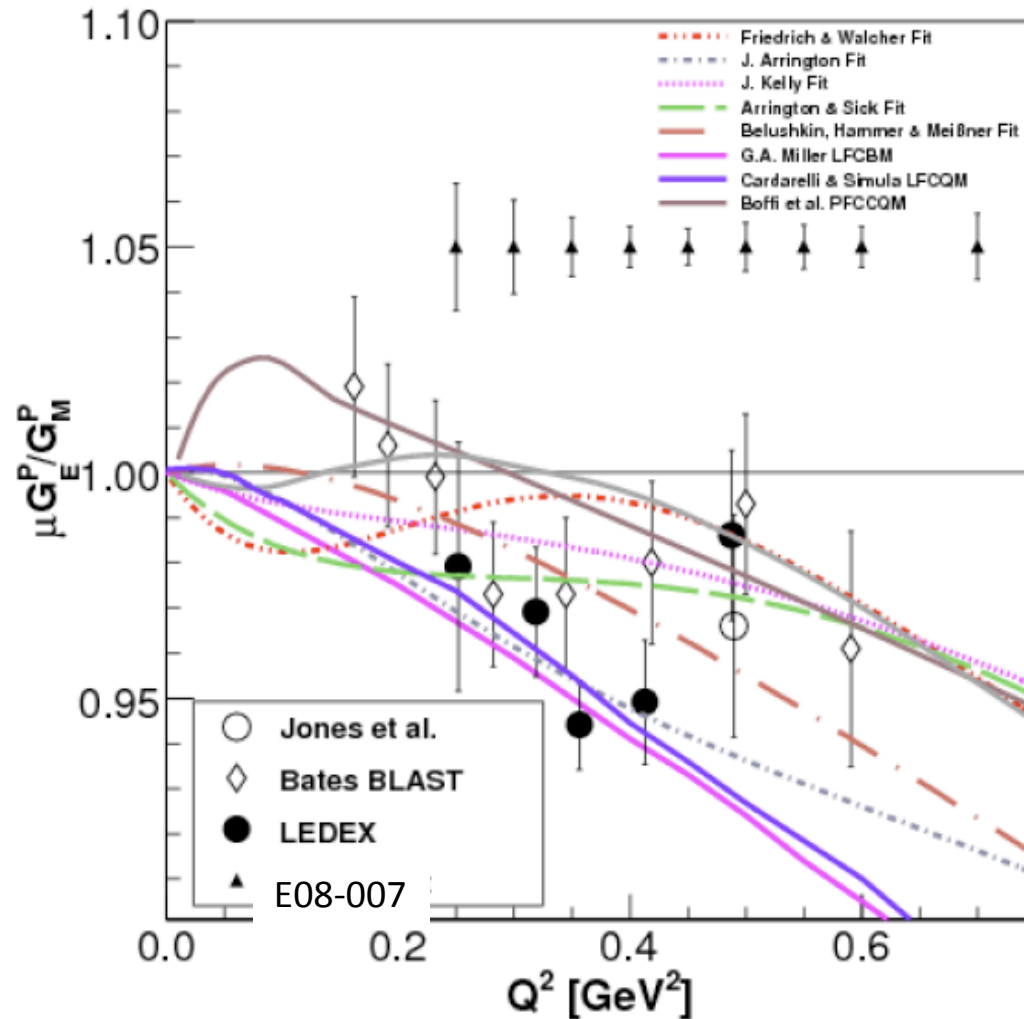
High precision proton form factor data at low Q^2

- 2003 - Fit by Friedrich & Walcher Eur. Phys. J. A17, 607 (2003):
- Smooth dipole form + "bump & dip" around $Q^2 \sim 0.1 \text{ GeV}^2$
- All four FFs exhibit similar structure at small momentum transfer.
- Proposed interpretation: effect of pion cloud peaked around $\sim 0.9 \text{ fm}$



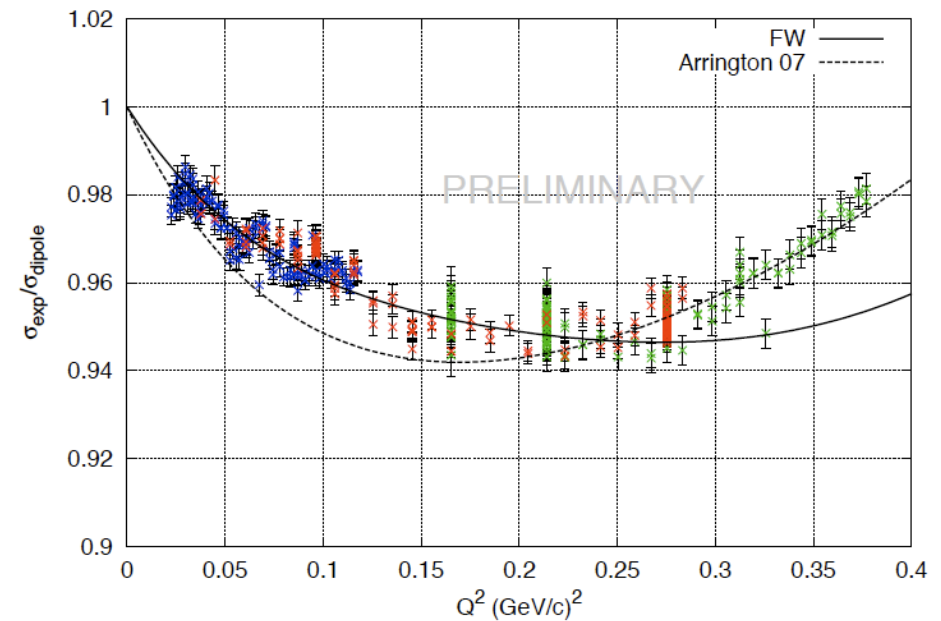
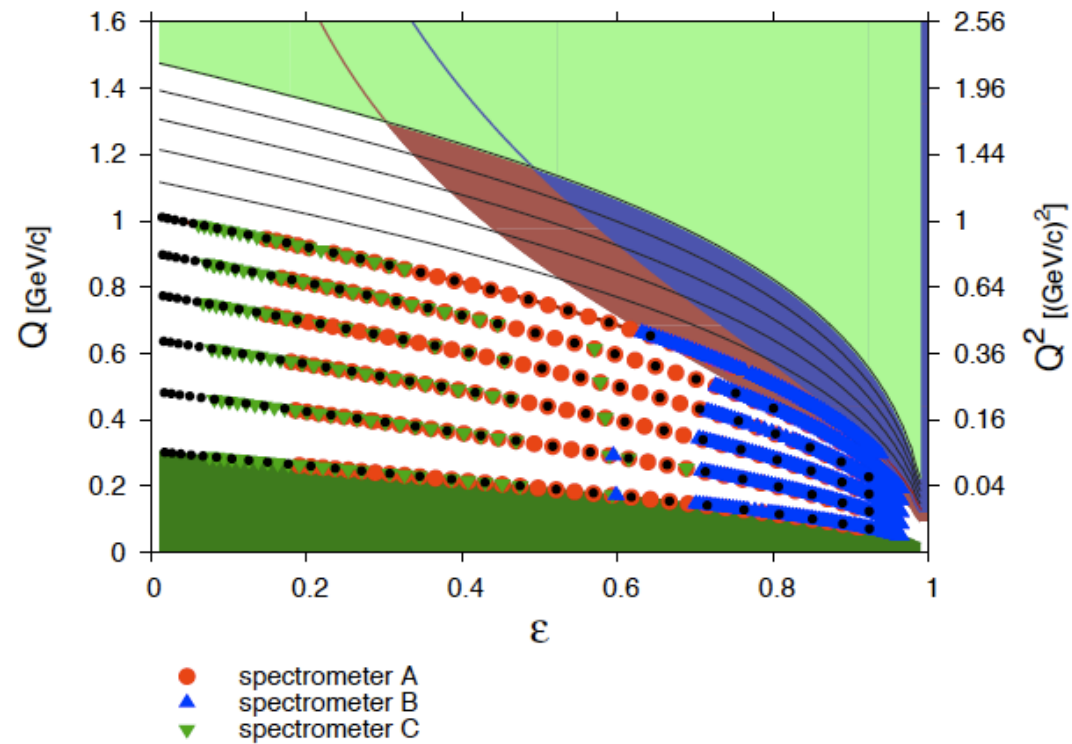
High precision proton form factor data at low Q^2

- BLAST and Hall A LEDEX result clearly show $\mu G_E^P/G_M^P < 1$ for $Q^2 \sim 0.3$ - 0.4 GeV^2 .
- But these does not support F&W analysis; different structure.
- New high precision experiment E08-007 in Hall A: results expected soon.



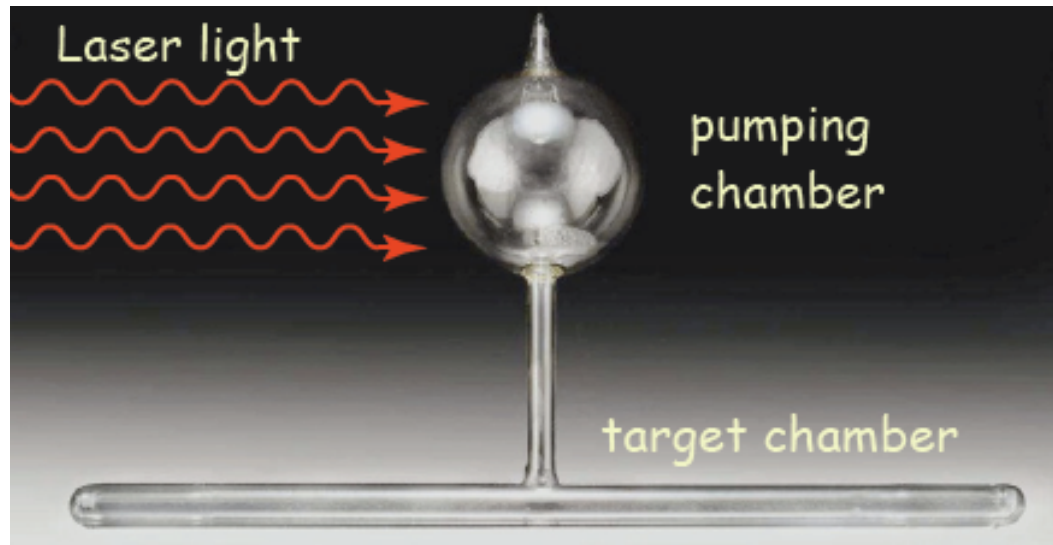
Mainz A1 high precision proton form factor experiment at low Q^2

- High precision Rosenbluth separation.
- ~ 1000 settings



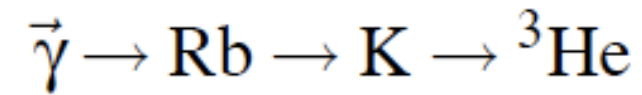
Summary

- Exciting and very active hadron form factor program thanks to high quality polarized beams, high-power polarized targets, large neutron detectors and high efficiency polarimeters.
- Electro-Magnetic form factors: high Q^2 precision data available
 - G_E^p : up to 8.5 GeV^2 ; discrepancy between polarization and Rosenbluth data may be due to 2-photon exchange.
 - G_E^n : up to 3.5 GeV^2
 - G_M^n : up to 5 GeV^2
- With 12 GeV upgrade these ranges can be extended to 15, 10 and 18 GeV^2 .
- Stringent tests on theoretical models.
- New high precision low Q^2 data: insight on Pion cloud ?
- Also a very active program measuring strange form-factors.

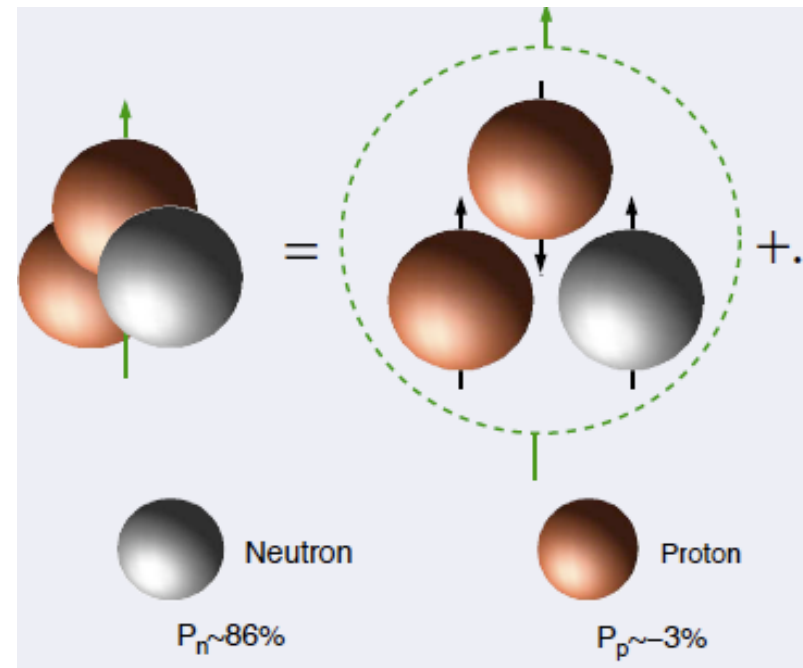
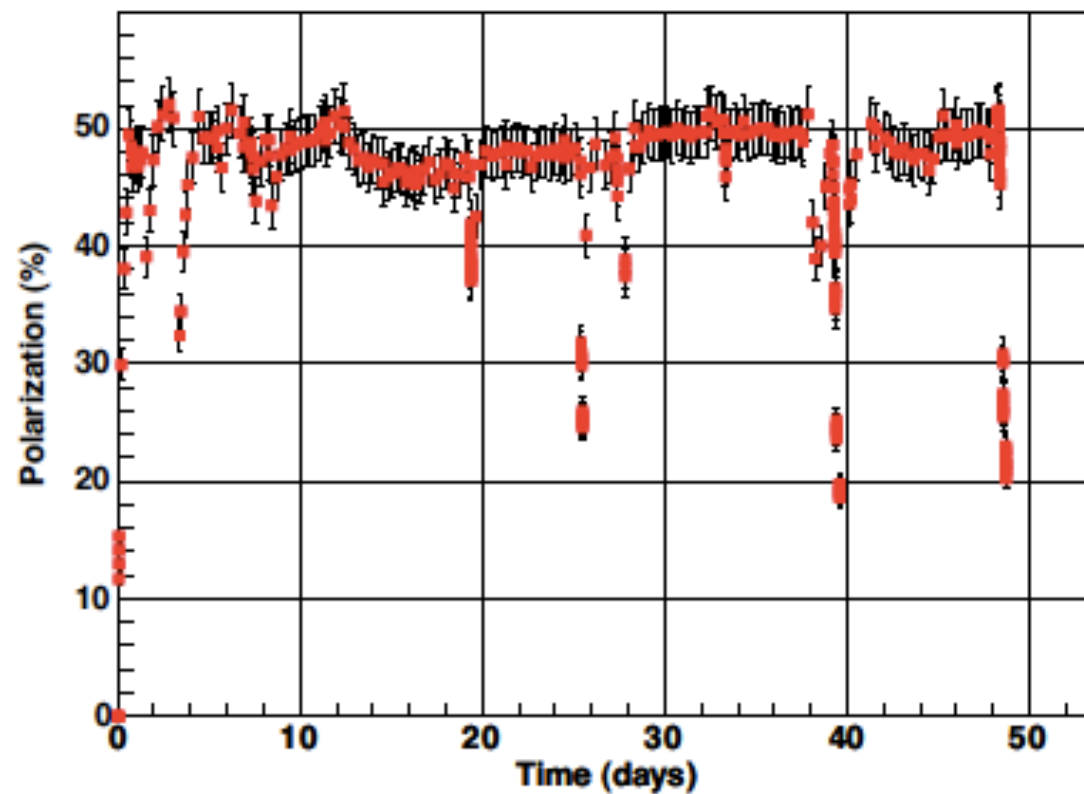


Polarized ^3He target

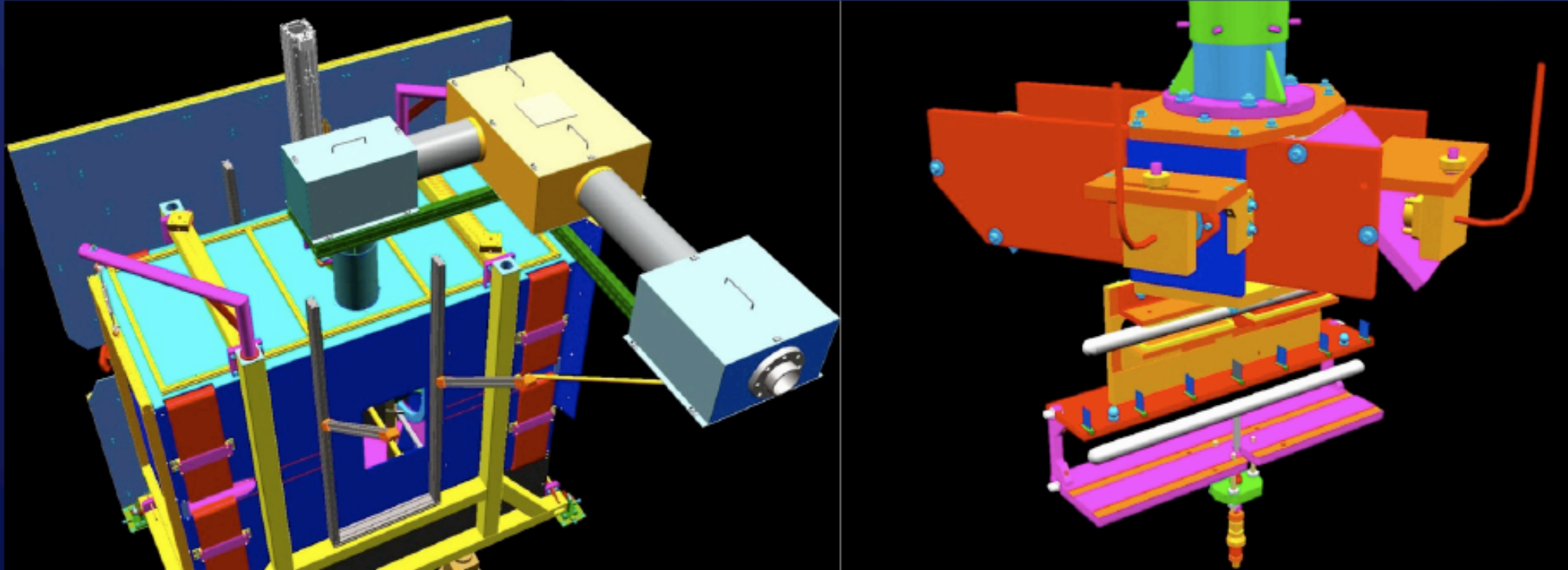
- based on spin exchange optical pumping
- Use of new Rb+K hybrid mixture to achieve higher spin exchange efficiency



- ~ 50% polarization achieved for most of running



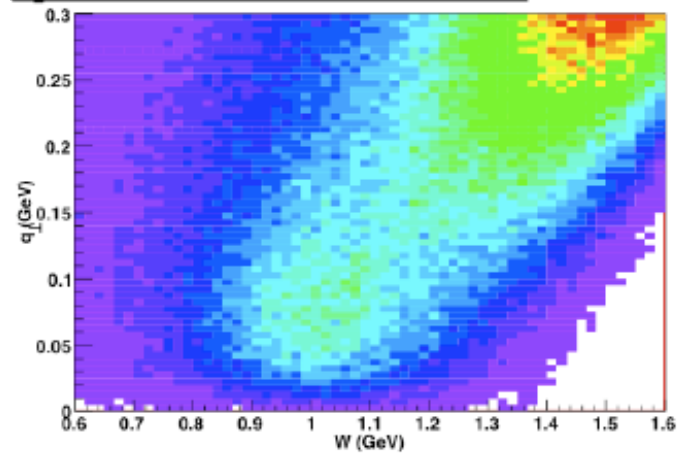
GEN required a novel implementation of ^3He spin-exchange technology



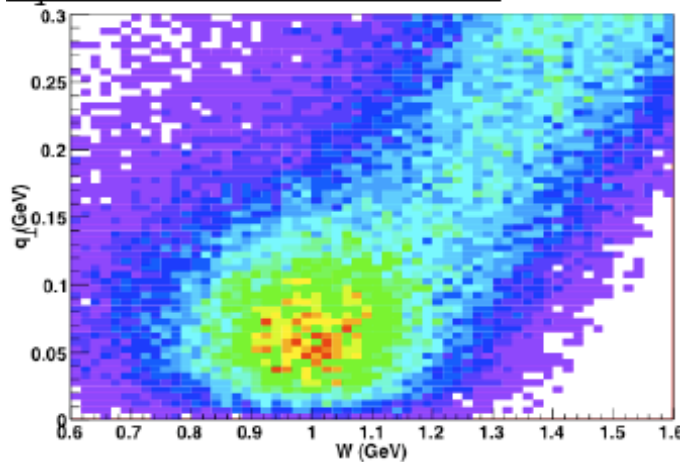
- Needed to be extremely close to the open-geometry BigBite magnet.
- Needed magnetic field inhomogeneities no worse than around 10 mG/cm
- "Iron Box" magnet design permitted target/BigBite distance of $\sim 1\text{m}$.
- Fiber-optic based laser/optics design greatly reduced space requirements.

Analysis

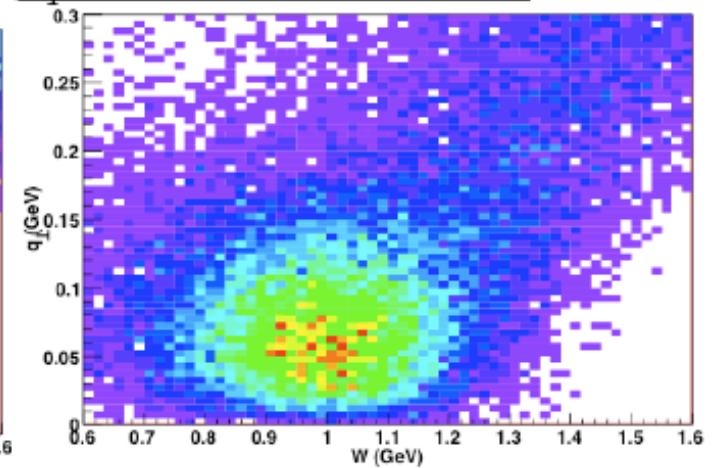
q vs. Invariant Mass, $Q^2 = 3.5 \text{ GeV}^2$, Raw spectrum



q vs. Invariant Mass, $Q^2 = 3.5 \text{ GeV}^2$, ToF cut



q vs. Invariant Mass, $Q^2 = 3.5 \text{ GeV}^2$, Full cuts



- Also a cut on missing mass of $3\text{He}(e,e',n)X$
- Use a MC simulation to evaluate inelastic contribution: very small
- Asymmetry corrected for:
 - p \rightarrow n conversion in shielding
 - Accidental events
 - A_{par} contribution and A_{proton}
 - Target and beam polarization
 - FSI for $3\text{He}(e,e',n)$ - not done yet.